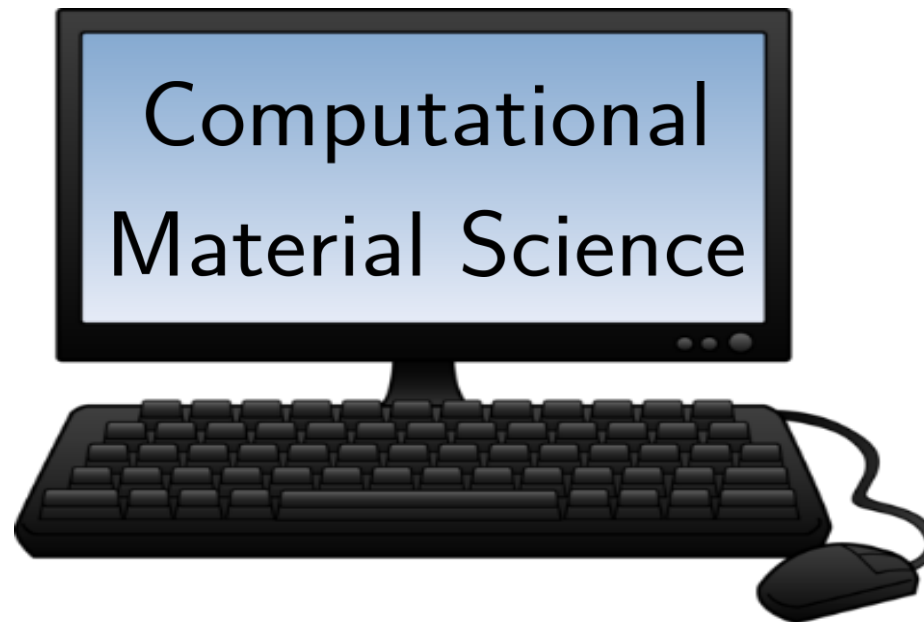
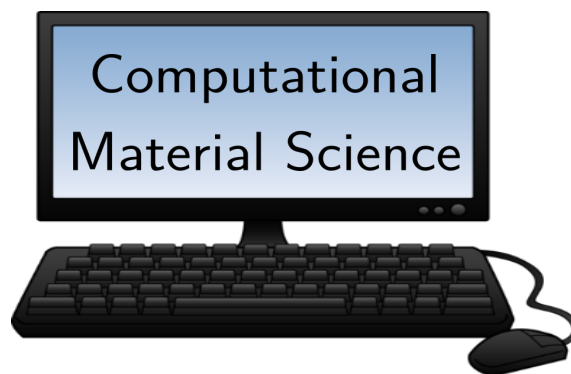


Selected Challenges in Material Simulations

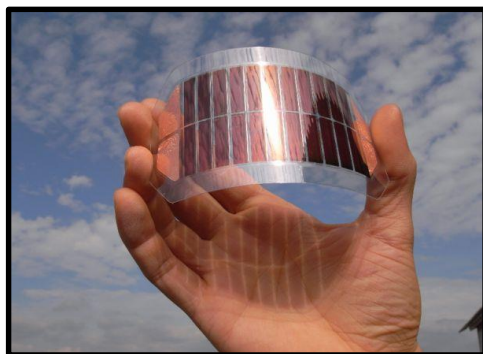


Oliver T. Hofmann

Graz University of Technology
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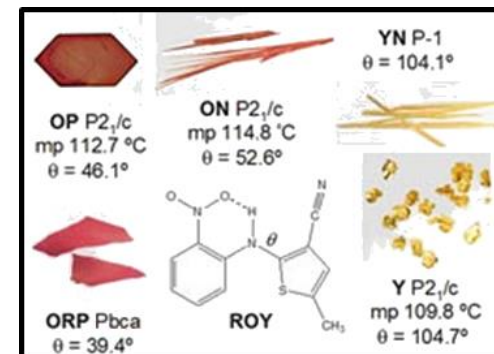
The quest for improved materials



Photovoltaic cells

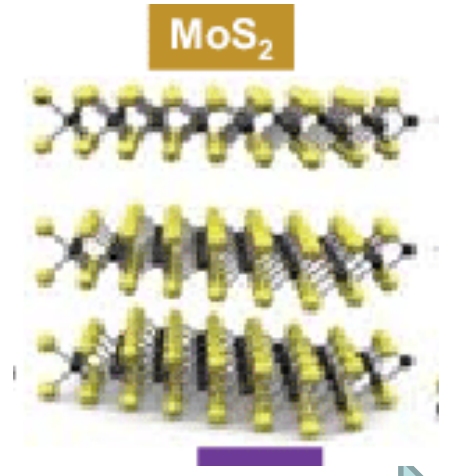
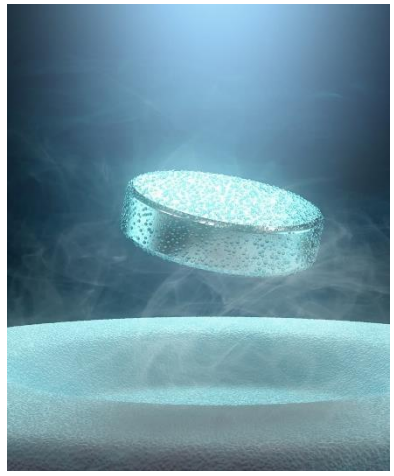
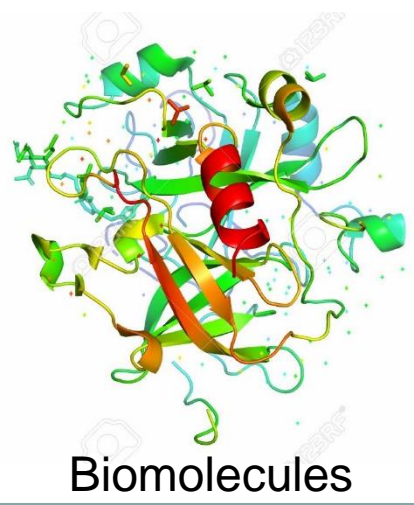


OLED displays

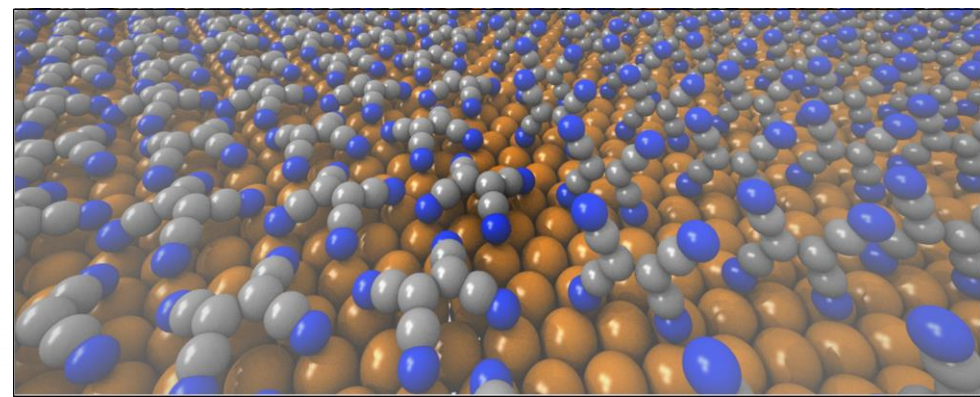
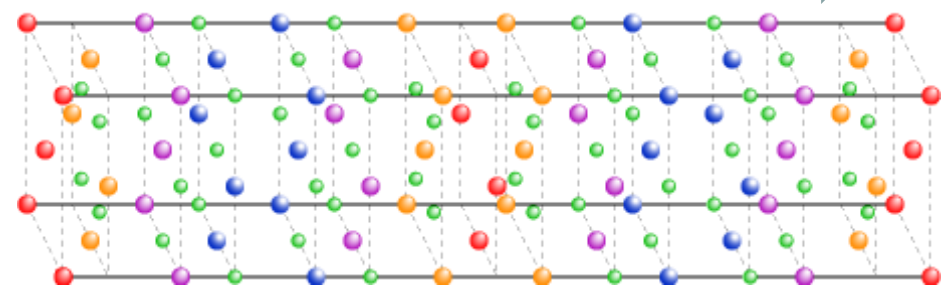
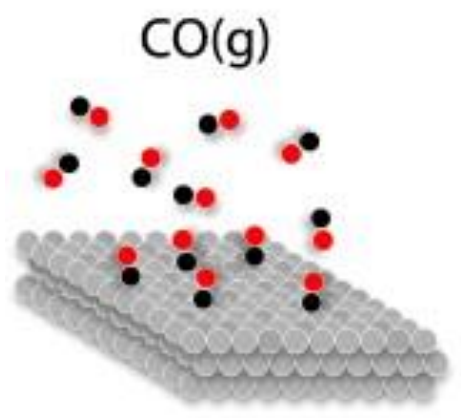


Pharmaceuticals

Si 14
28,9
Silicium



Transition
metal
oxides



- **Magnetic materials / excitations**
- **Finite-temperature effects / Anharmonicities**
- **Polymorphism**
- **Charge Transfer**

- Magnetic materials / excitations
- Finite-temperature effects / Anharmonicities
- Polymorphism
- Charge Transfer



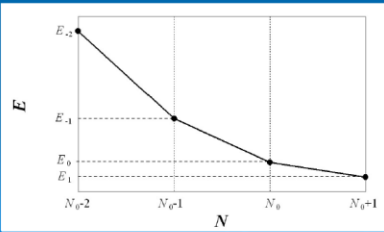
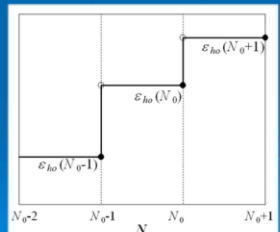
**Piecewise linearity in DFT:
Exact results from ensemble arguments**

Fractional no. of electrons: $N = N_0 + \alpha$

- Piecewise-linearity
- The IP theorem

$E(N) = (1 - \alpha)E_0 + \alpha E_1$

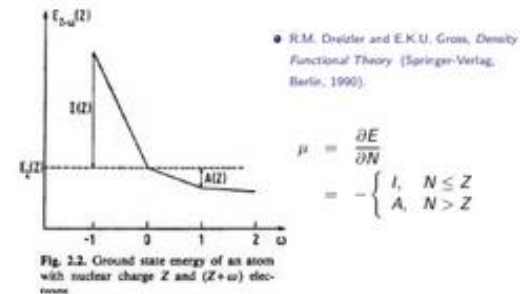
$\epsilon_{ho}(N) = -I(N_0) =: E_0 - E_{-1}$

Perdew, Parr, Levy, Balduz, PRL 49, 1691 (1982); Yang, Zhang, Ayers, PRL 84, 5172 (2000).
Janak, PRB 18, 7165 (1978); Levy, Perdew, Sahni, PRA 30, 2745 (1984).

Derivative discontinuity

As a function of N , the energy is a sequence of straight line segments

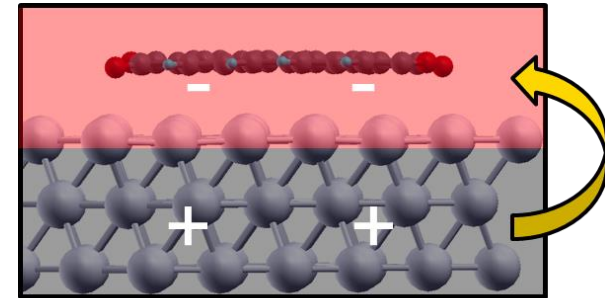


• R.M. Dreiter and E.K.U. Gross, Density Functional Theory (Springer-Verlag, Berlin, 1990).

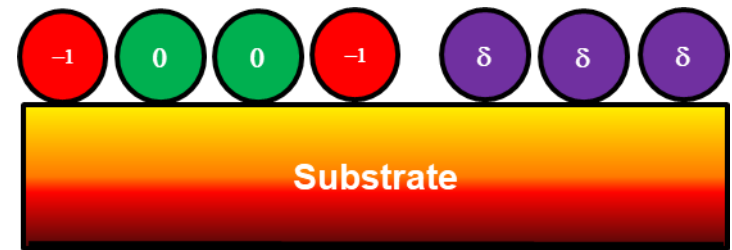


Outline: Issues with Charge Transfer

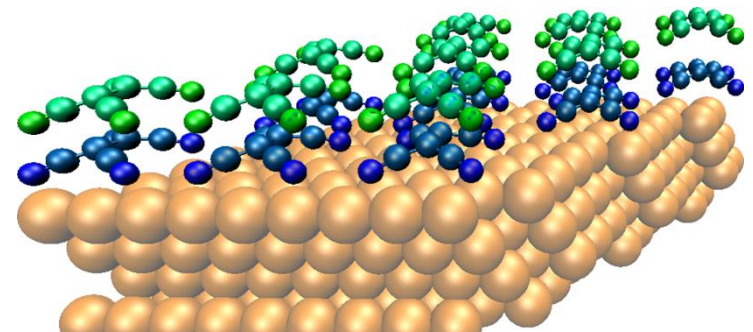
- Amount of charge transfer



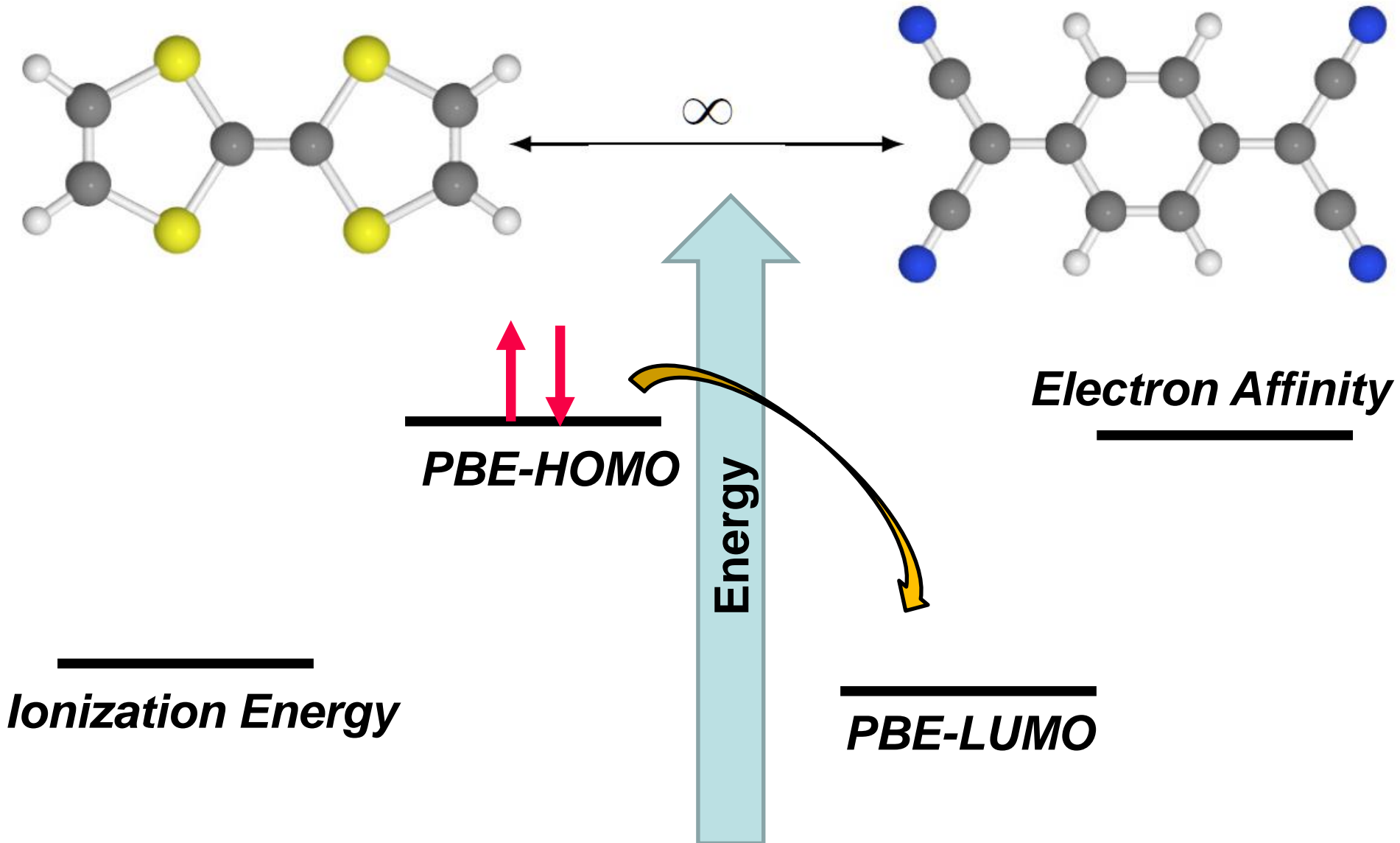
- Charge distribution



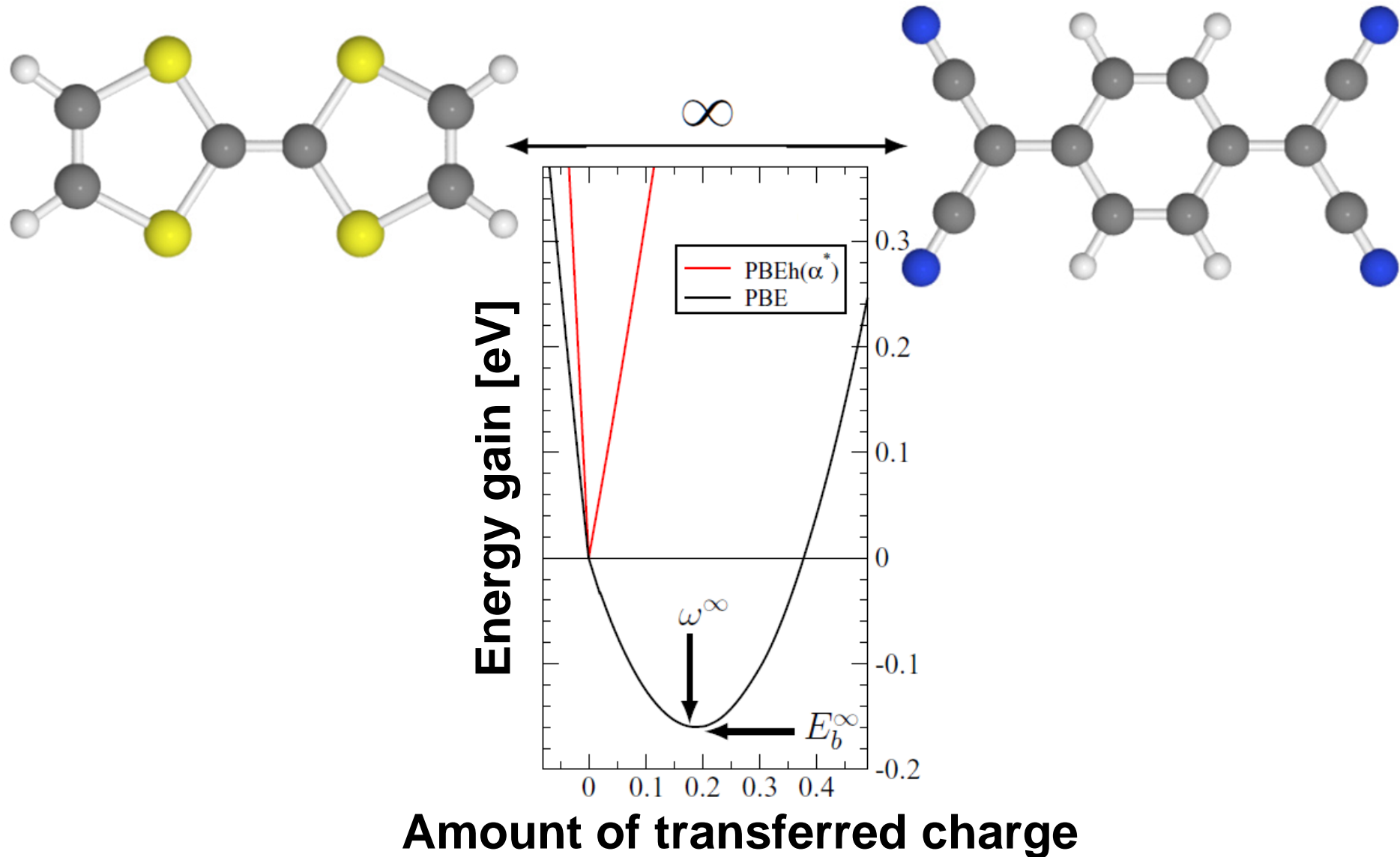
- Charge transfer \leftrightarrow Structure?

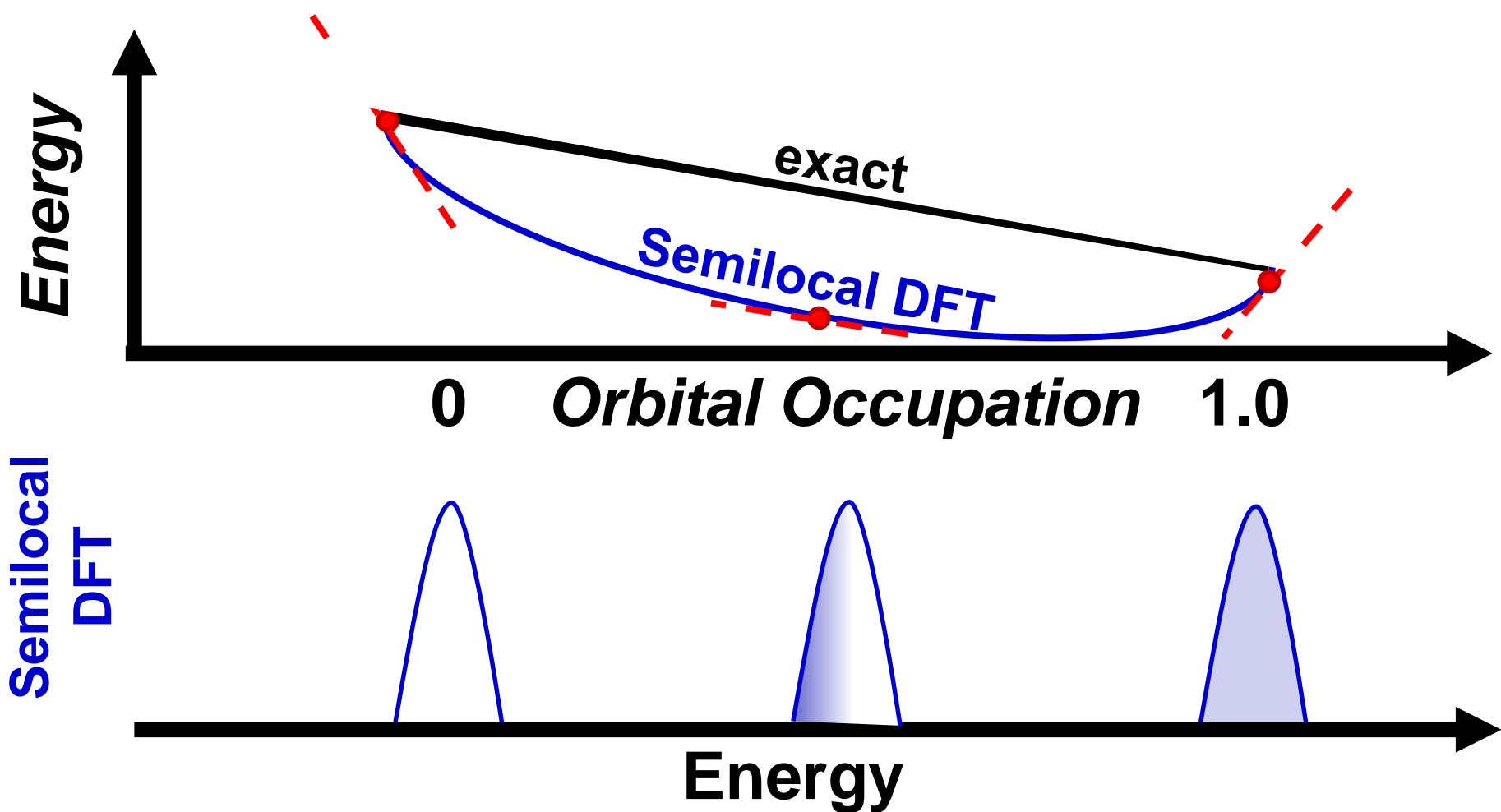


Charge Transfer In DFT: The Band Gap Problem

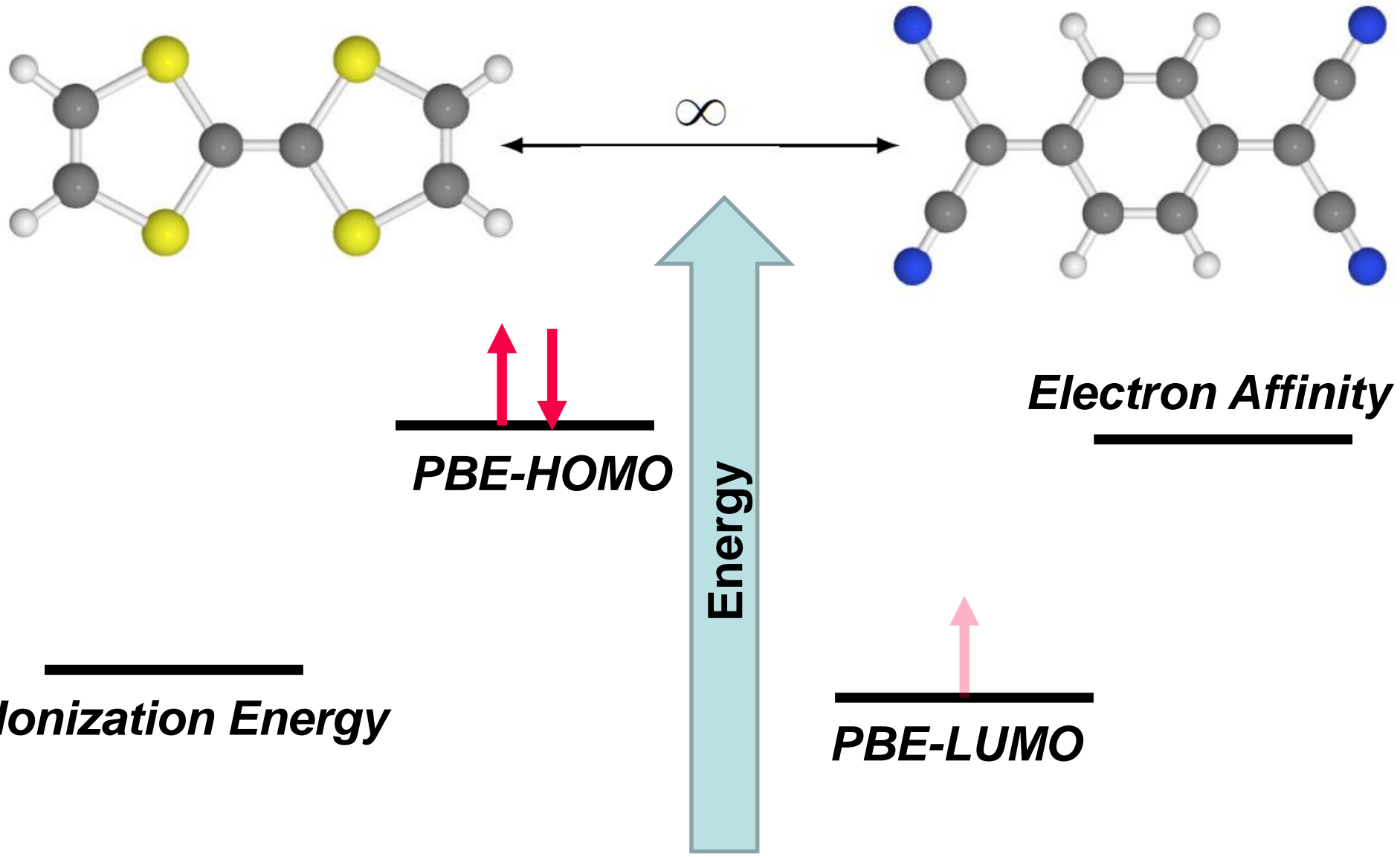


Charge Transfer In DFT: The Band Gap Problem

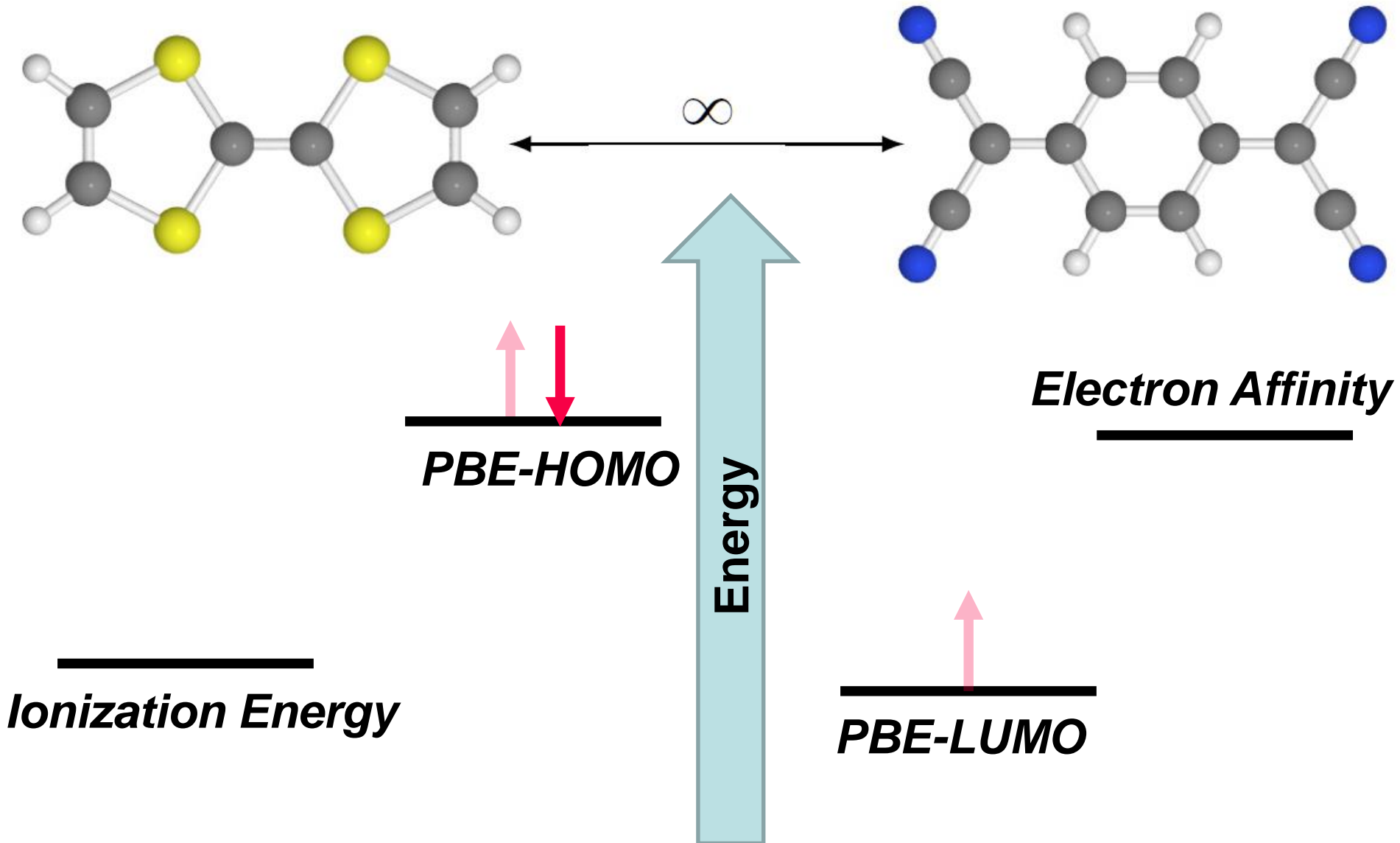


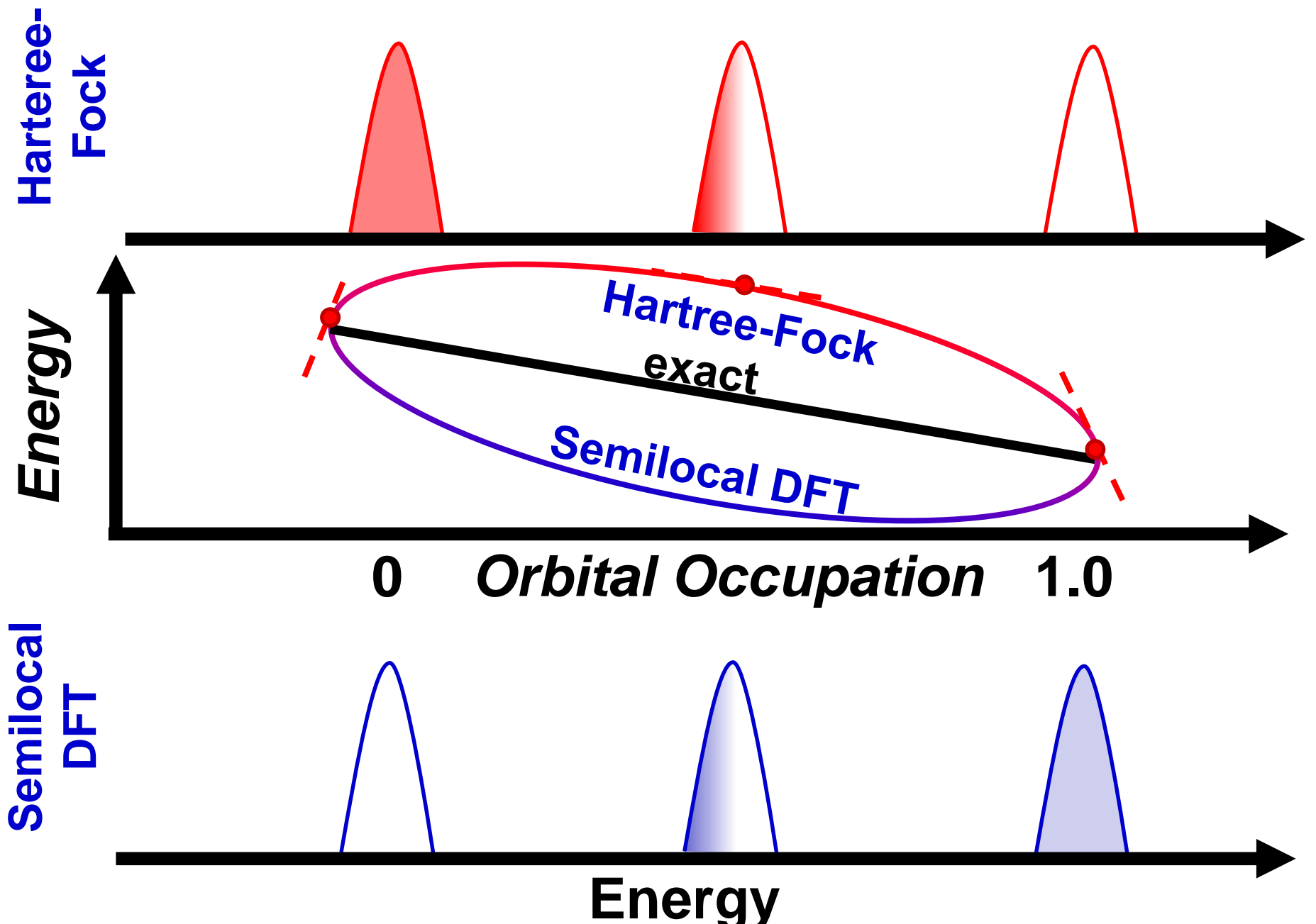


Charge Transfer In DFT: The Band Gap Problem

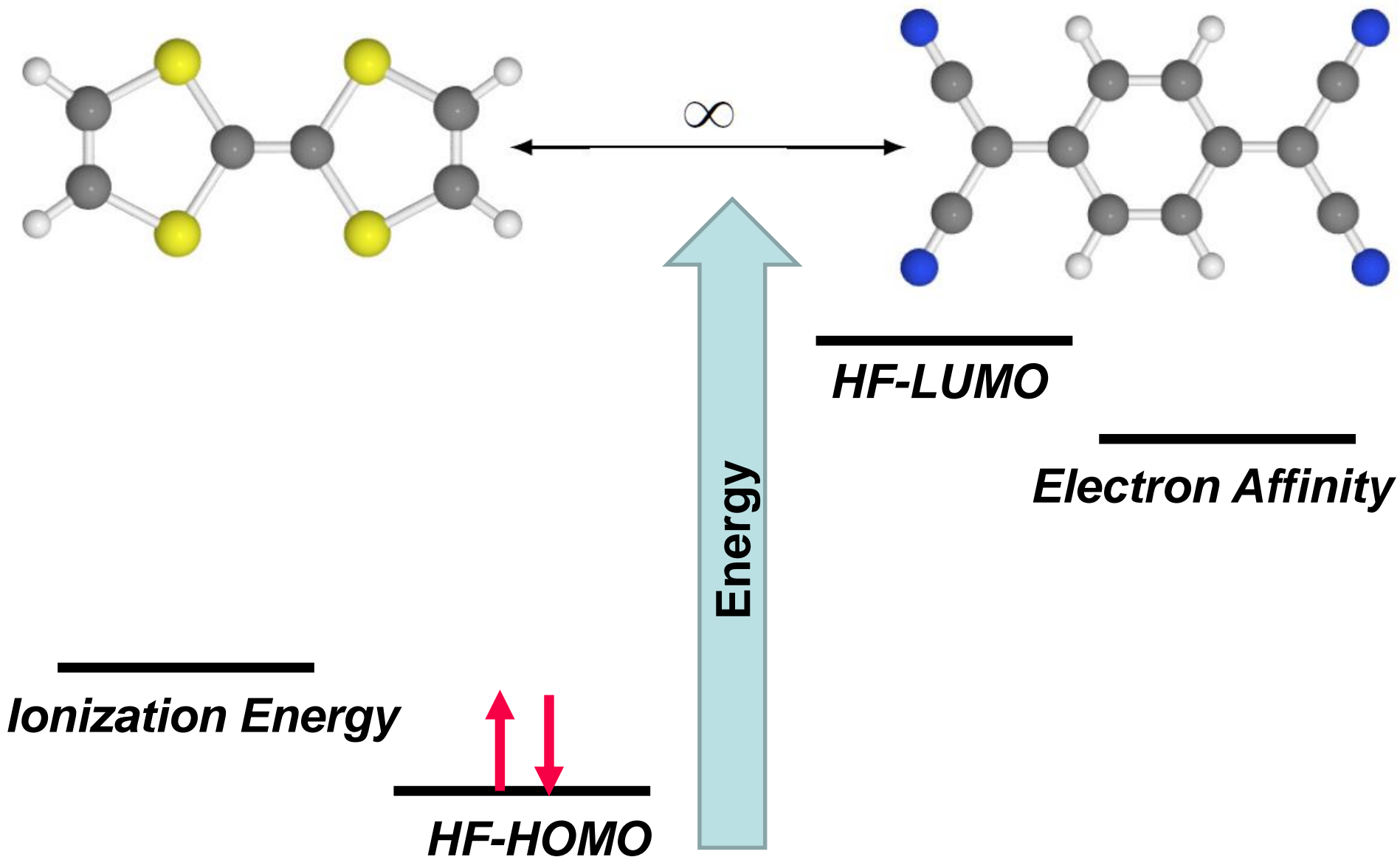


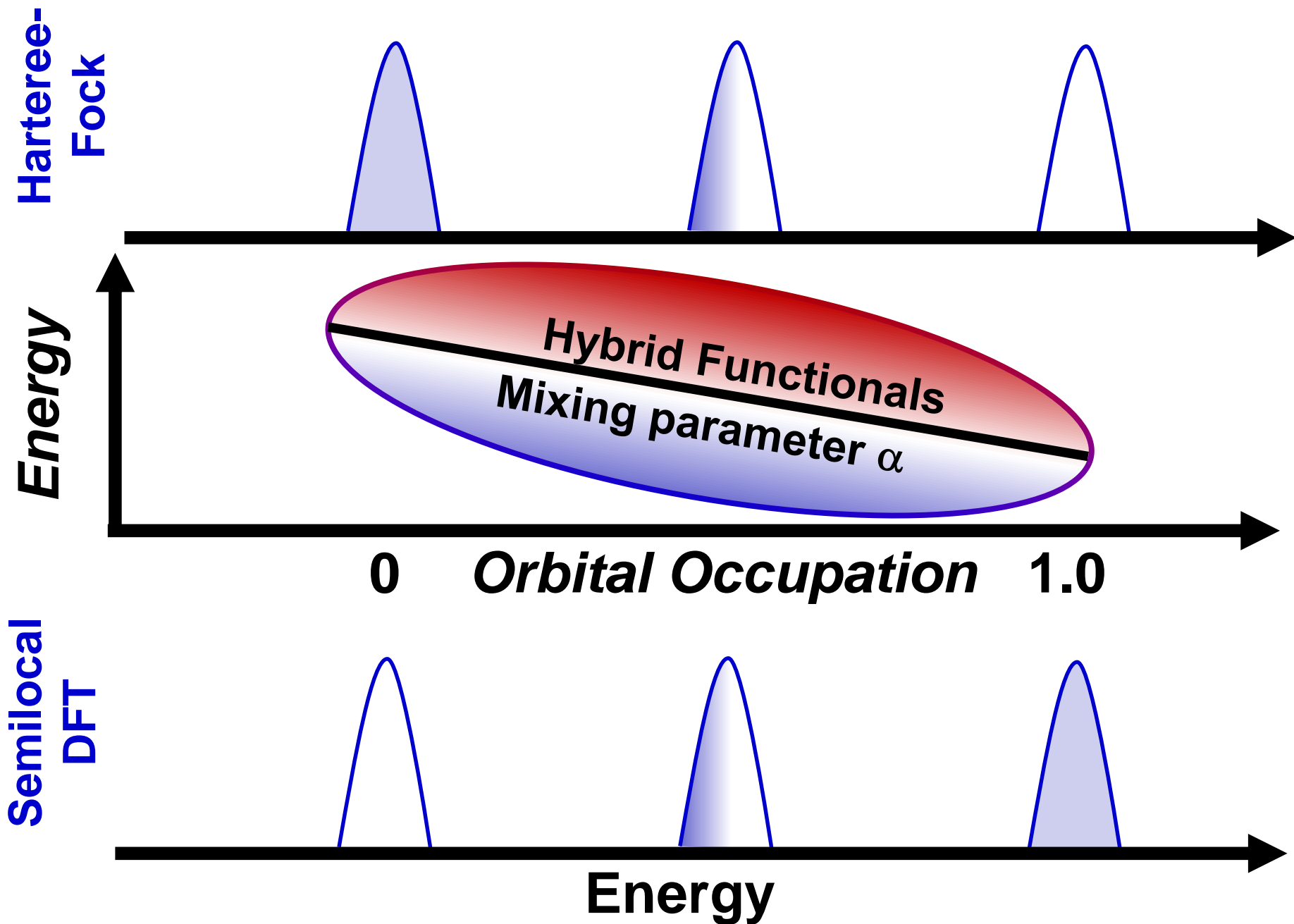
Charge Transfer In DFT: The Band Gap Problem

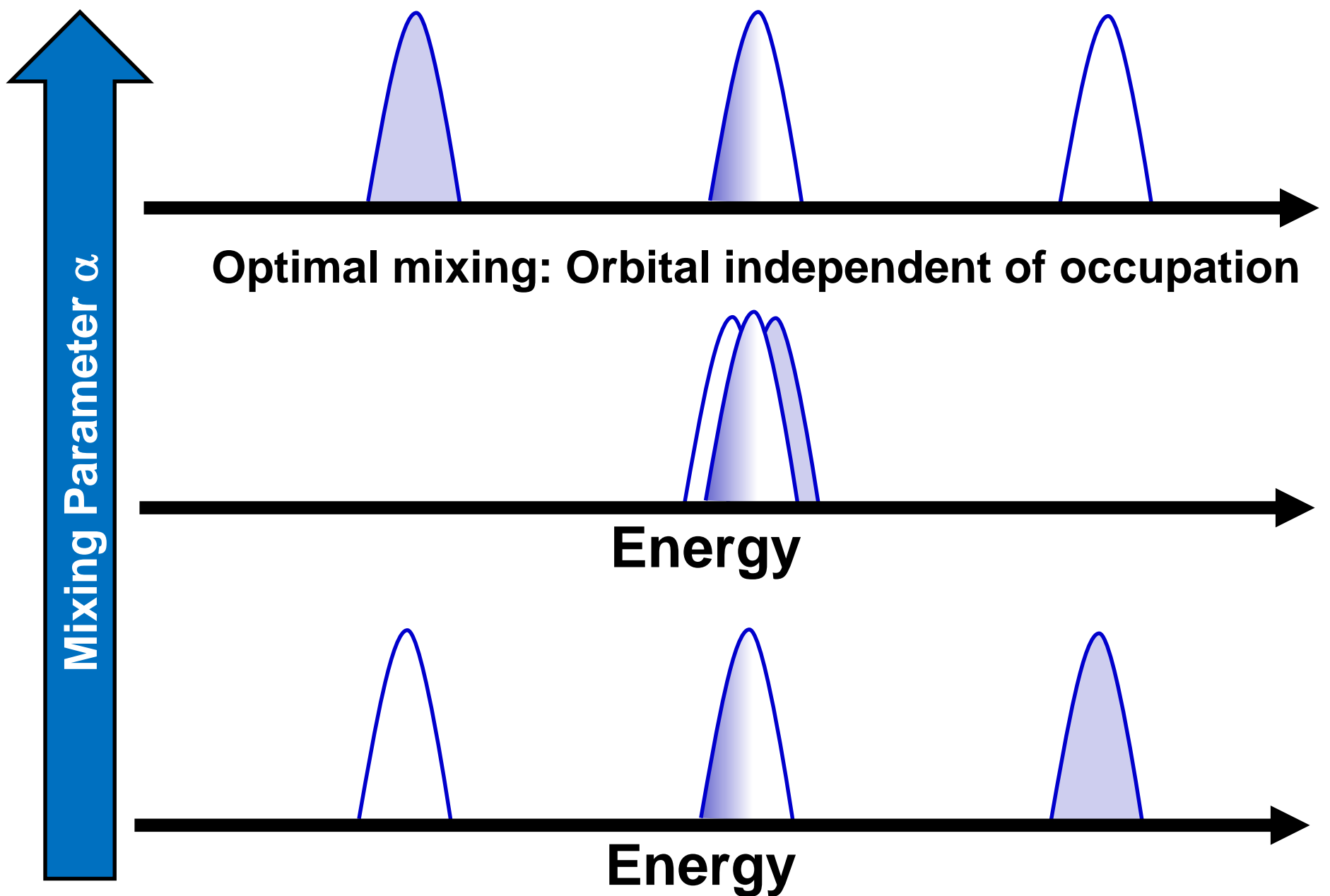




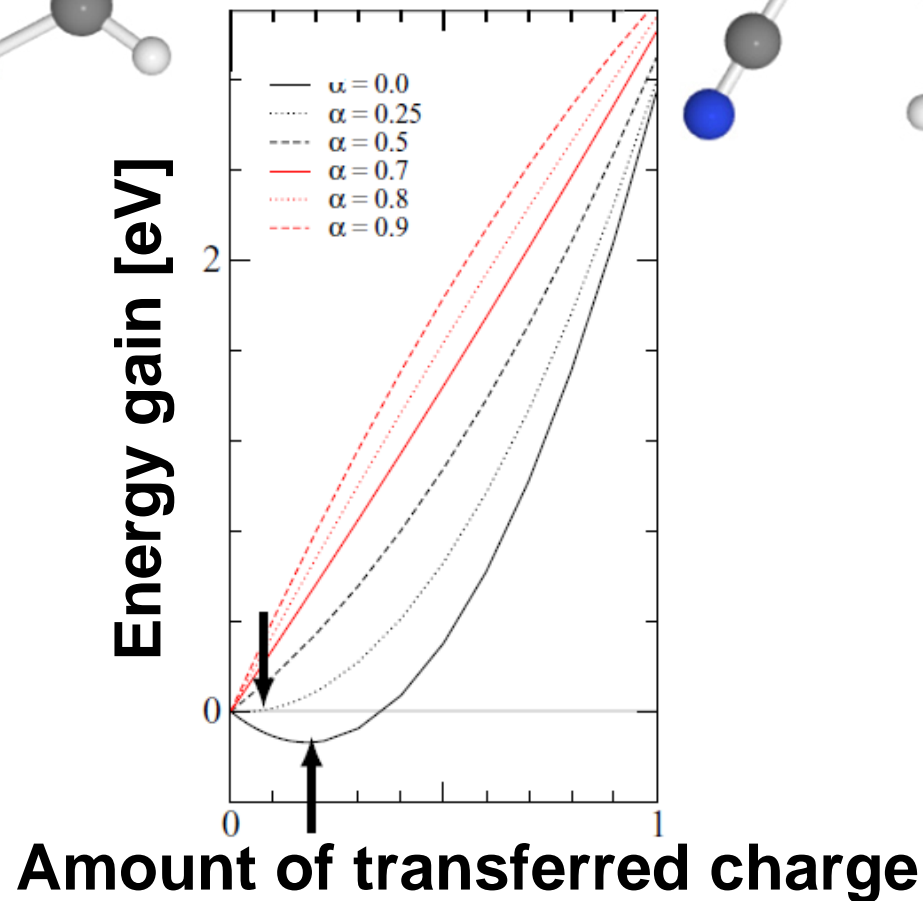
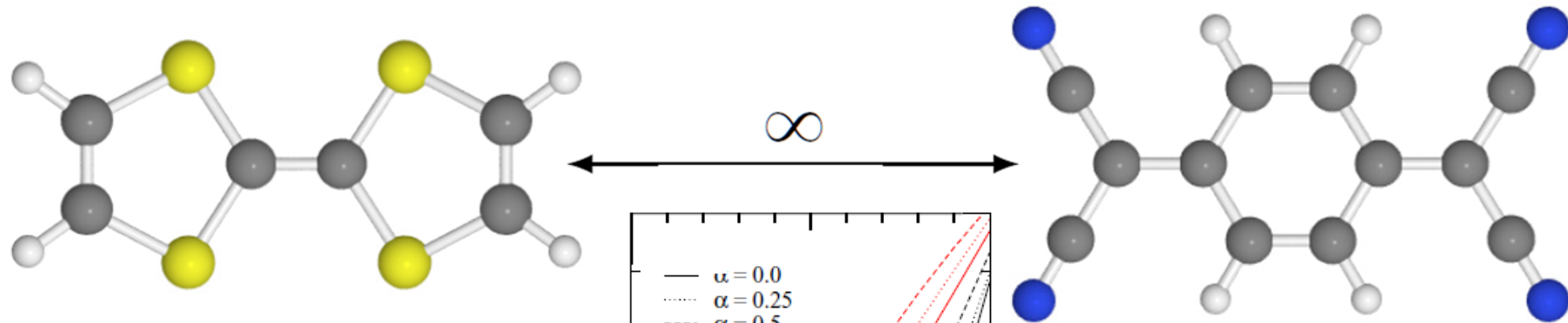
Charge Transfer In DFT: The Band Gap Problem



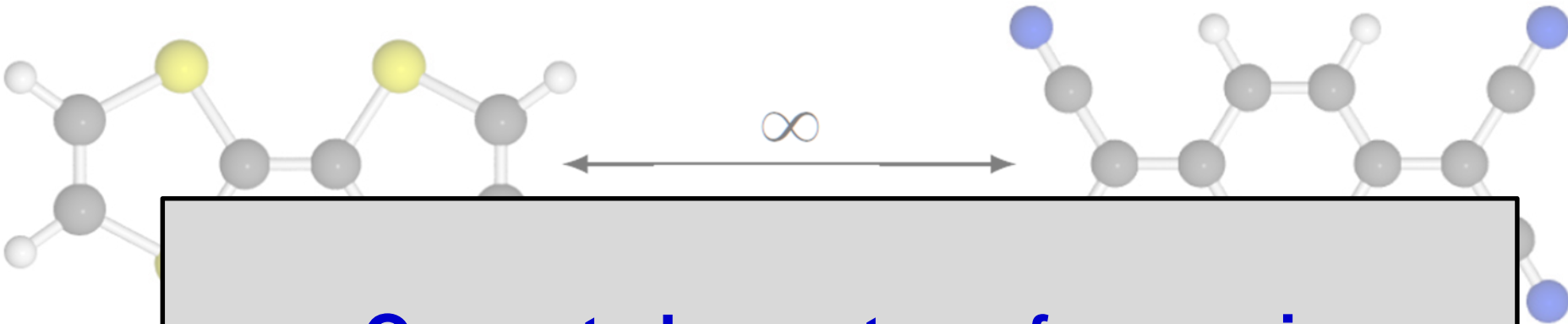




Charge Transfer In DFT: The Band Gap Problem

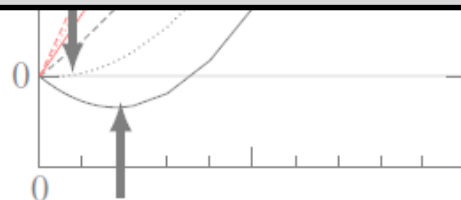


Charge Transfer In DFT: The band gap problem



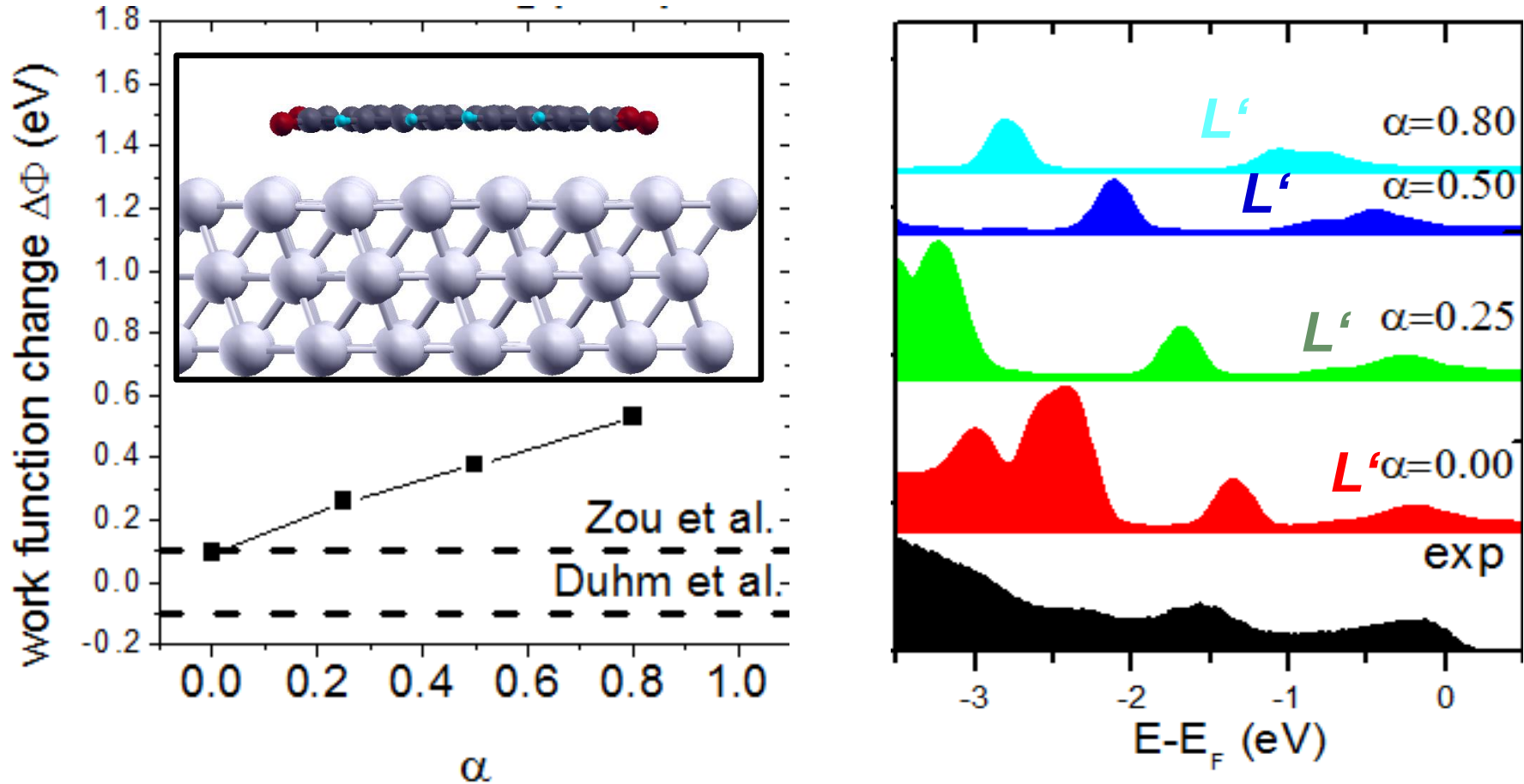
- **Correct charge transfer requires „optimal“ mixing**
- **Hybrid functional decrease charge-transfer**

Karolewski et al., The Journal of Chemical Physics, 138 (20), 2013



Amount of transferred charge

Exceptions to the rule: Large CT



Filled PBE-LUMO gets pushed further below the Fermi-energy \rightarrow charge-transfer increases

The performance of semilocal DFT for CT:

➤ **Overestimate CT when it is small**

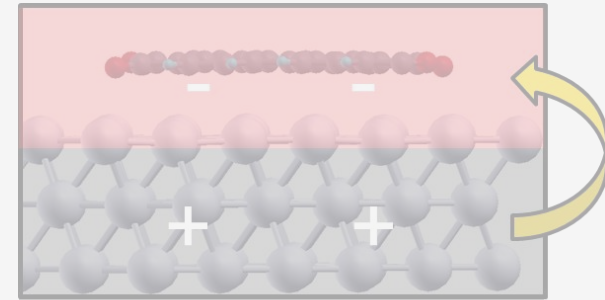
Atalla et al., Phys. Rev. B 94, 035140

➤ **Underestimate CT when it is large**

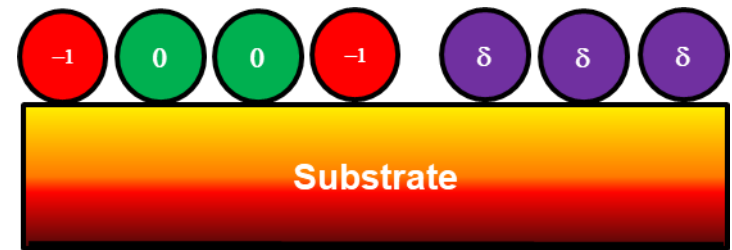
Hofmann et al., NJP, 15 (12), 123028, 2013

Outline: Issues with Charge Transfer

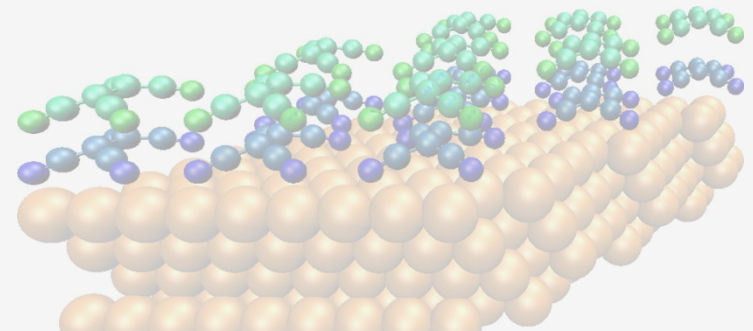
- Amount of charge transfer



- Charge distribution**

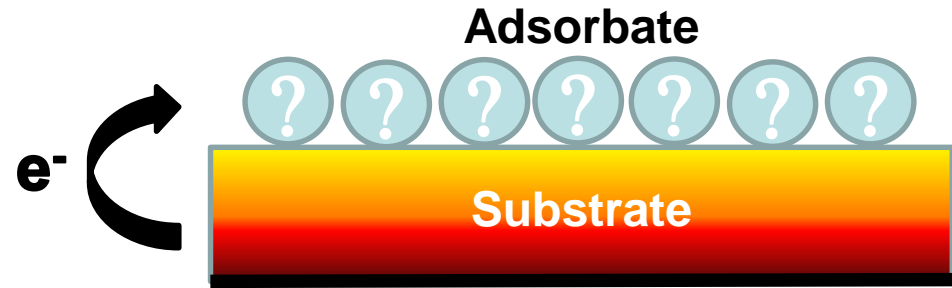


- Charge transfer \leftrightarrow Structure?





Braun et al., *Advanced Materials*, **2009**, (21), 1450



Experimental CT usually less than 1e / molecule

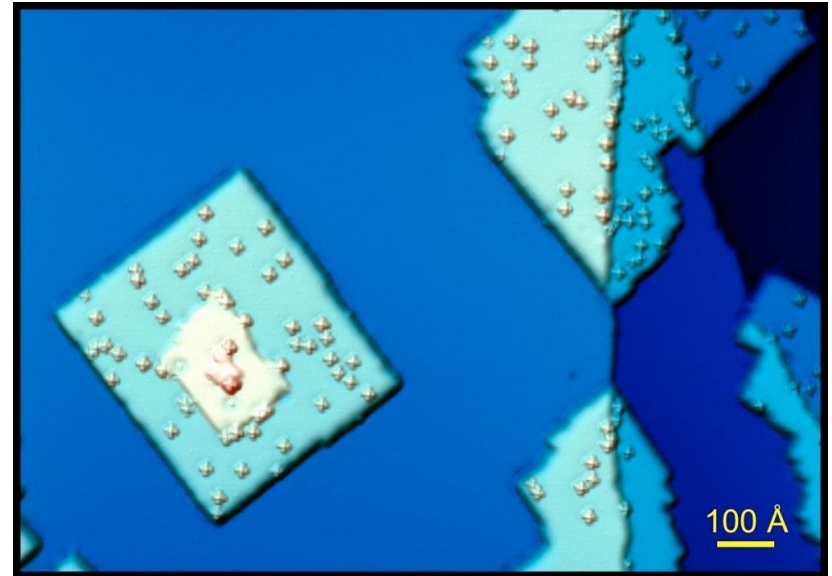
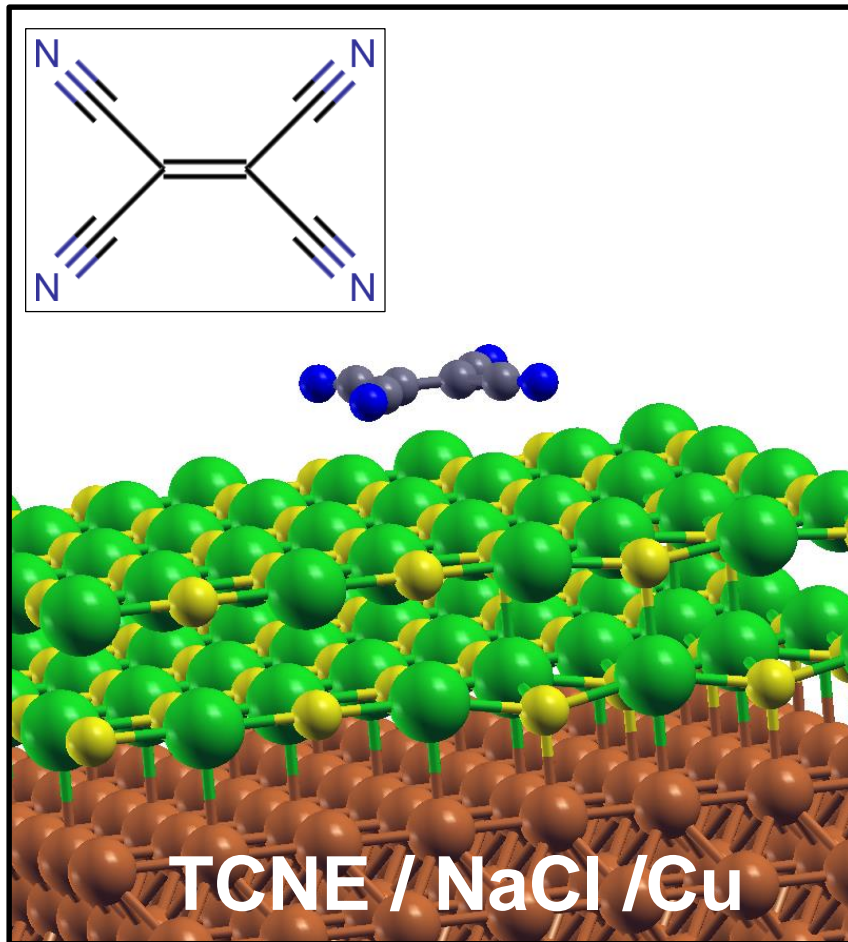
Integer Charge Transfer (ICT)

Not captured
semilocal functionals
(incl. +U)

Gruenewald et al., *J. Phys. Chem. C*, **2015**, 119 (9), 4865

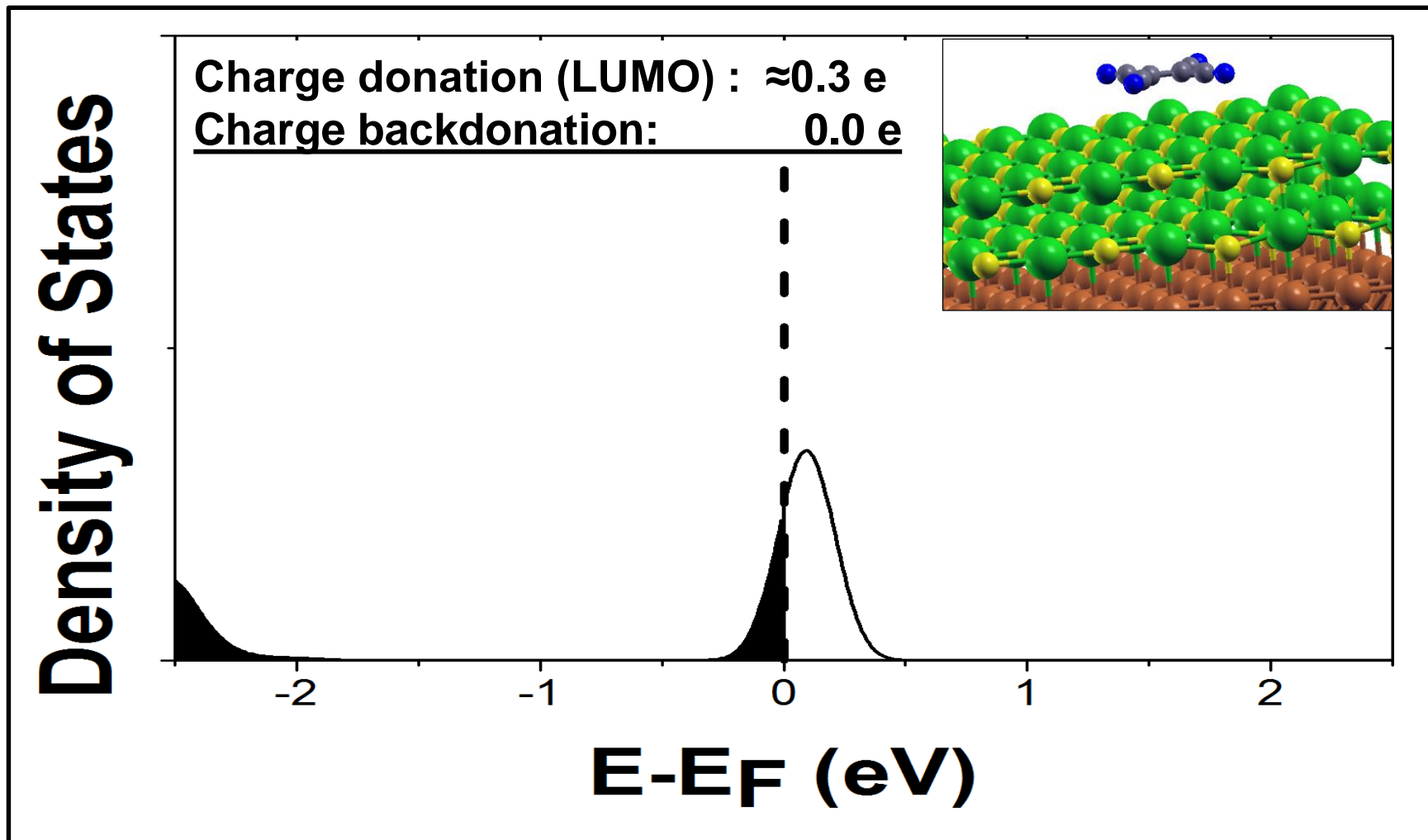
Fractional Charge Transfer (FCT)

Weak Interaction– TCNE/NaCl/Cu

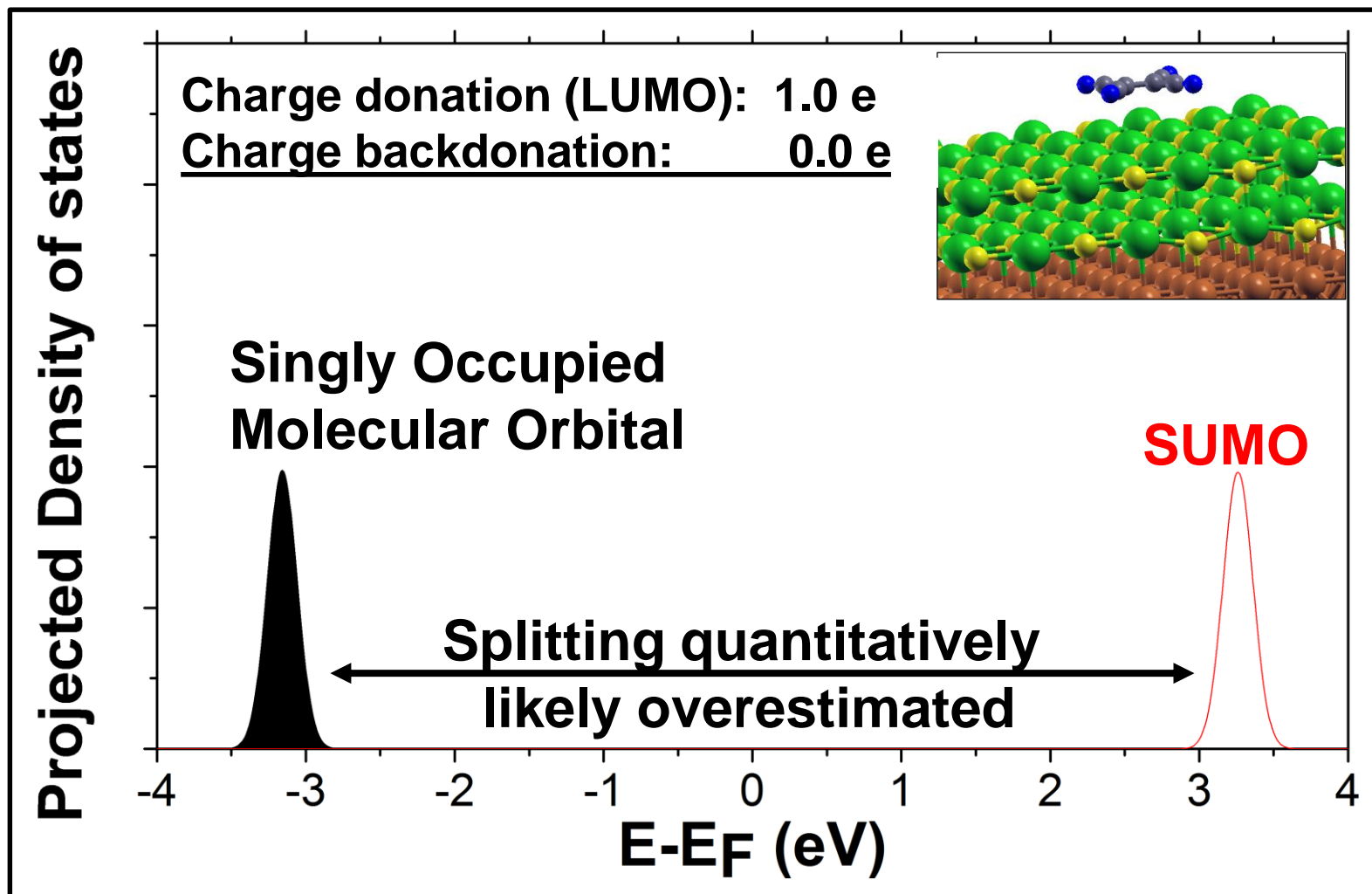


TCNE / NaCl / Ag, curtesy Daniel Wegner

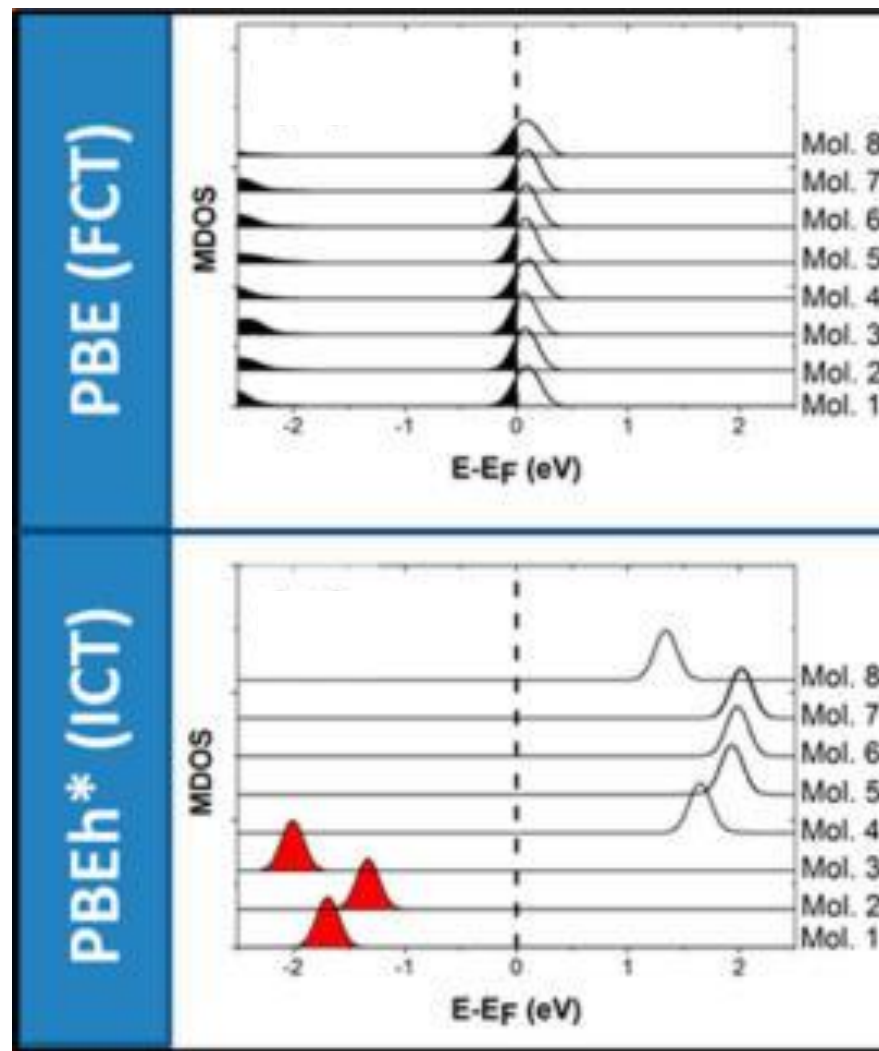
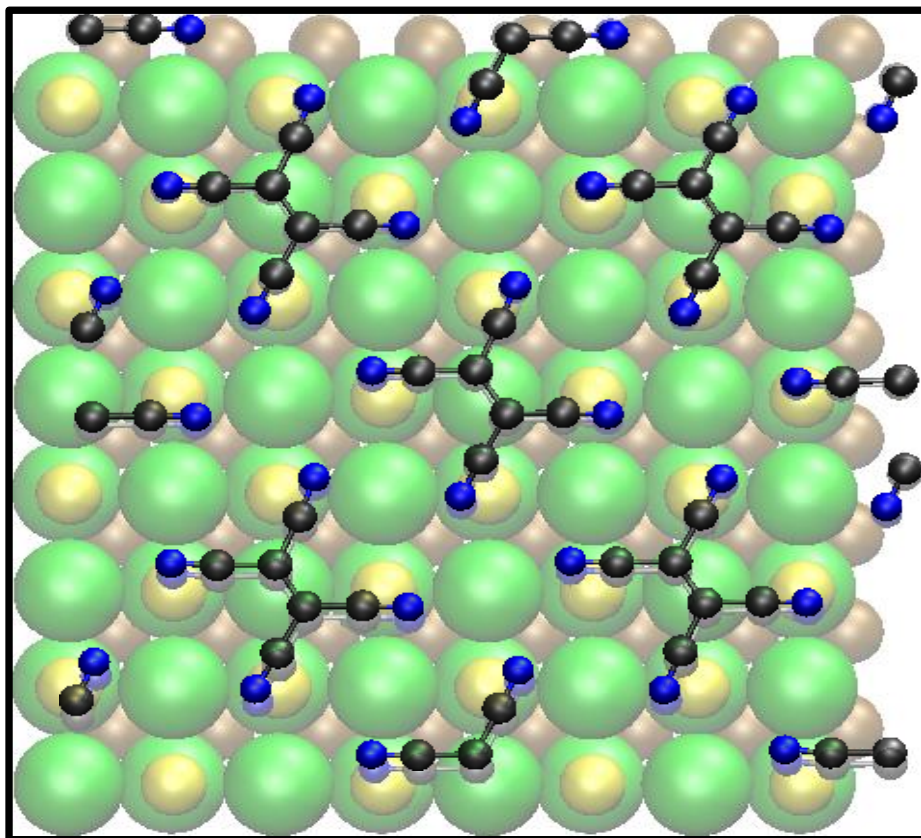
TCNE/NaCl/Cu – PBE calculation



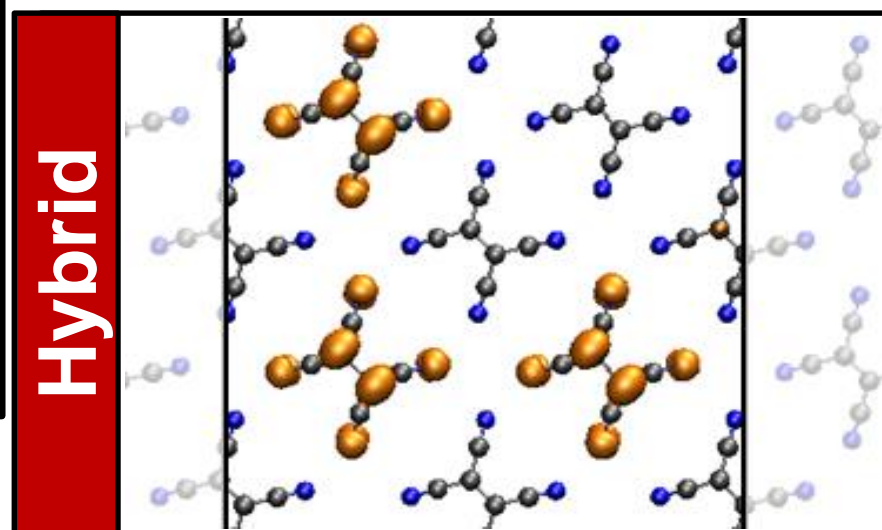
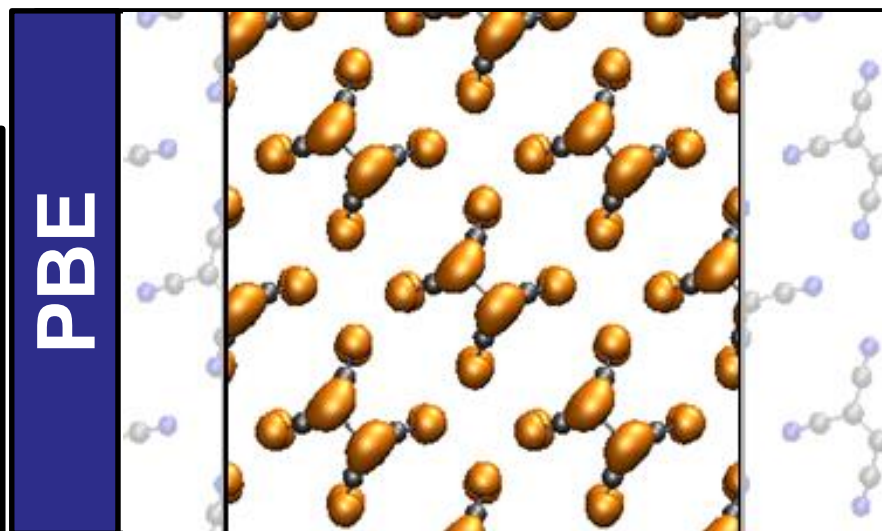
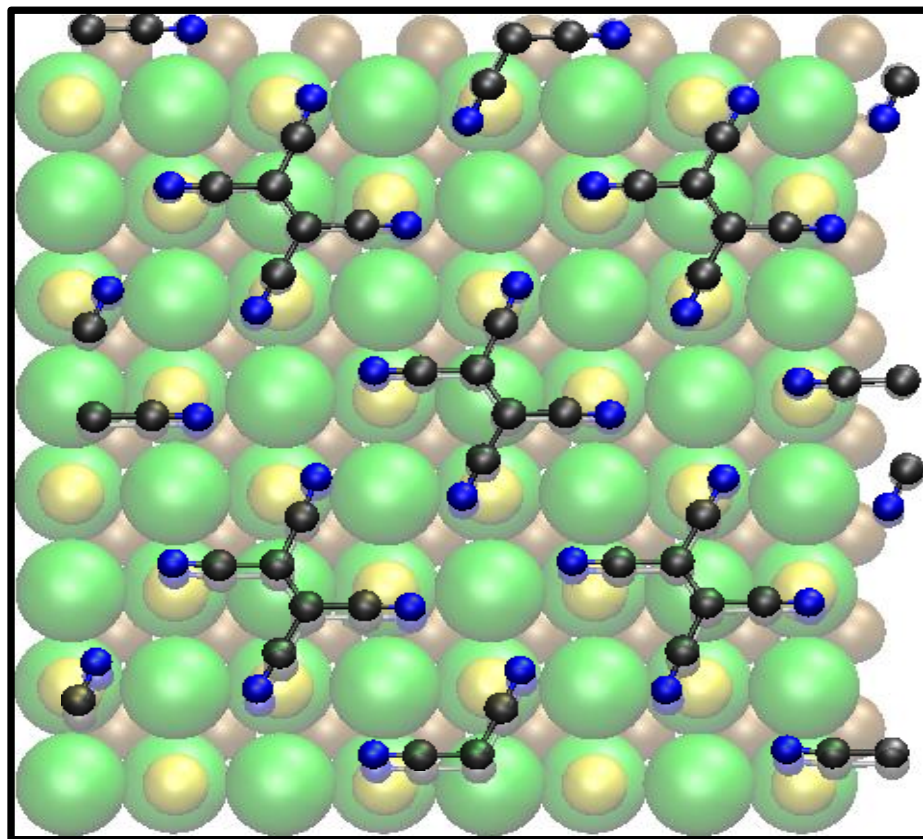
PBEh* ($\alpha = 0.8$) results – single molecule



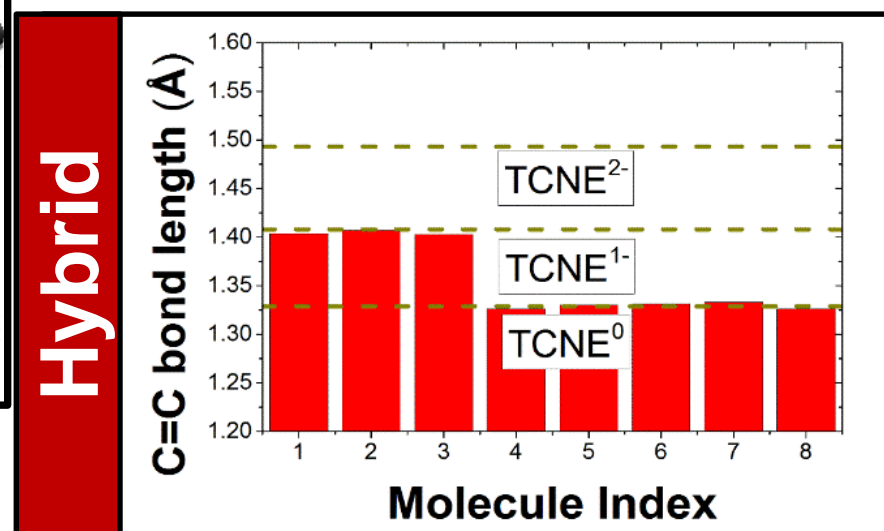
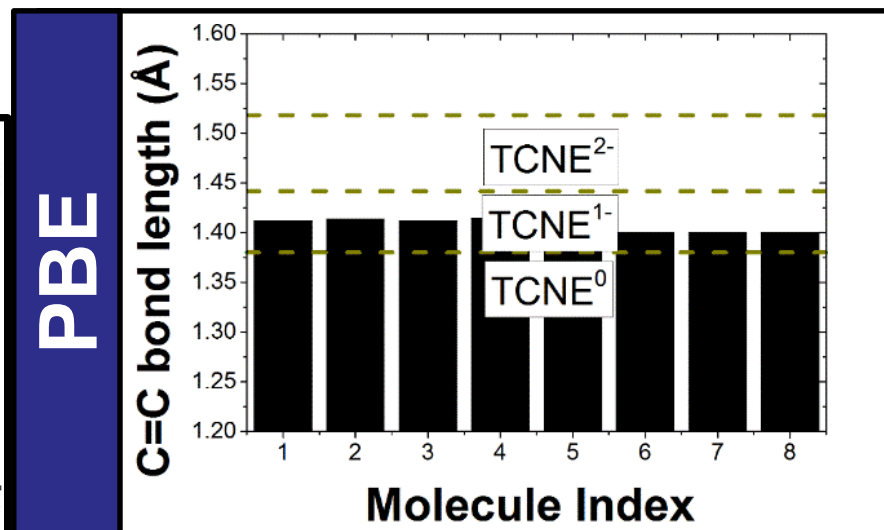
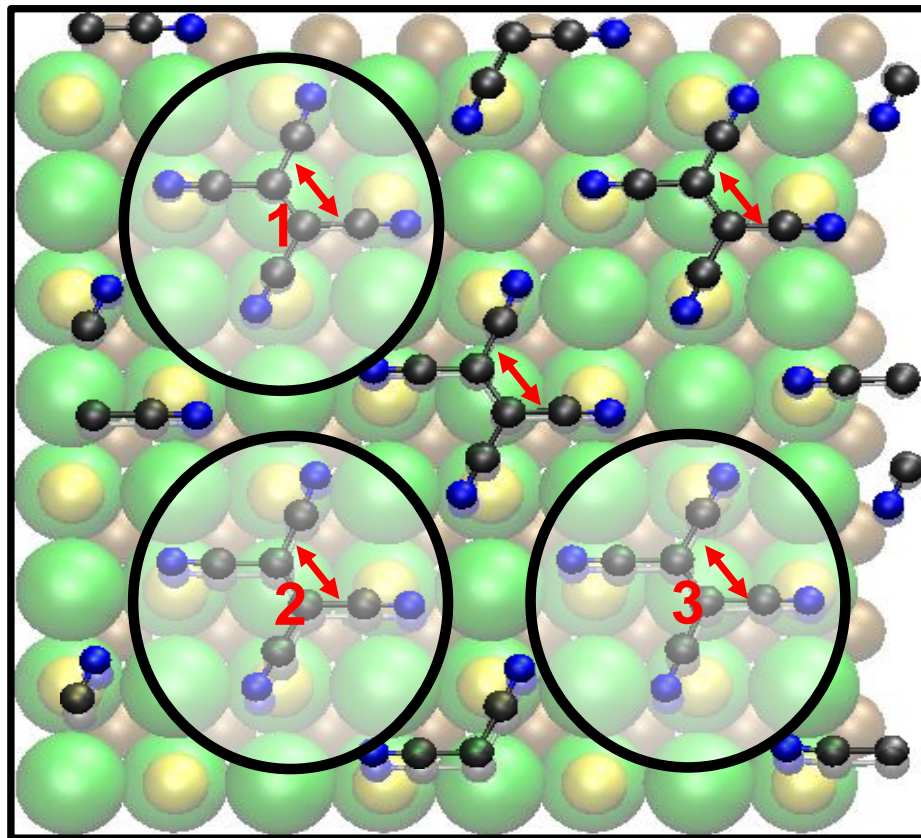
Density of States



Charge Distribution ...

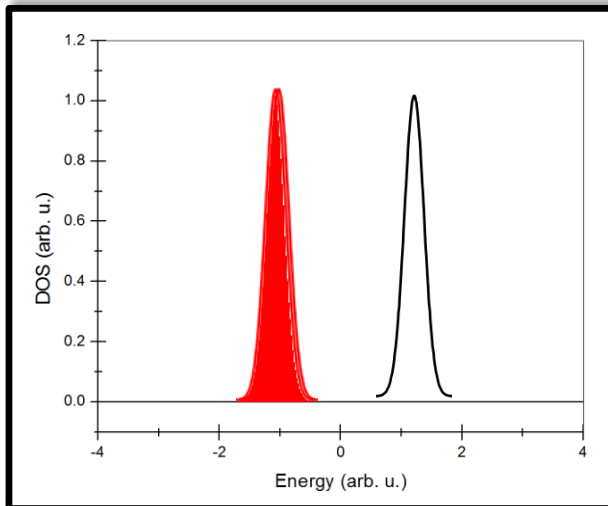


... and observables



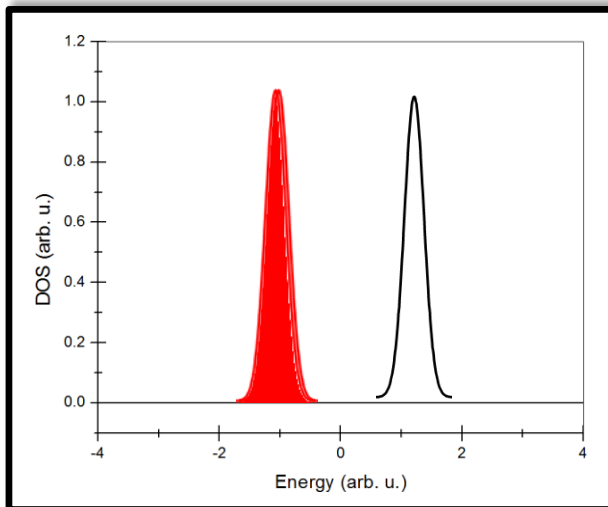
- ❖ Consider two empty states localized on two different molecules
- ❖ Symmetry broken (e.g., lattice distortion, numerical inaccuracy, etc.)
- ❖ Now fill with single charge

$$\frac{\delta\epsilon}{\delta n} > 0 \rightarrow \text{self-repulsion}$$

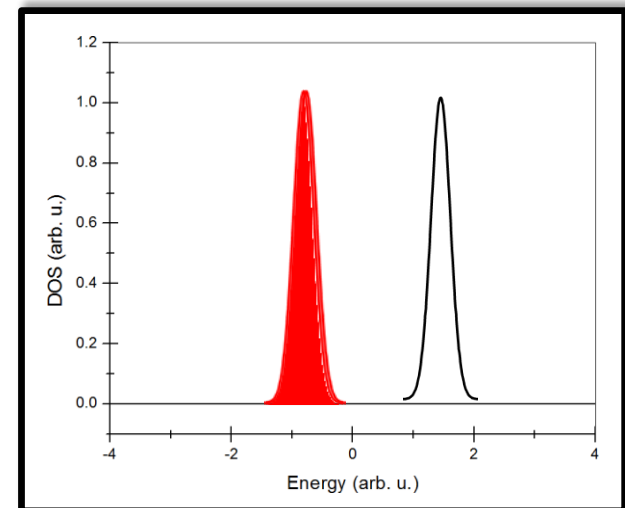


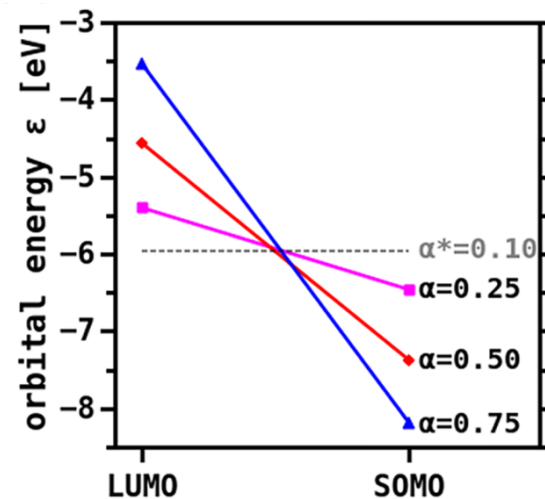
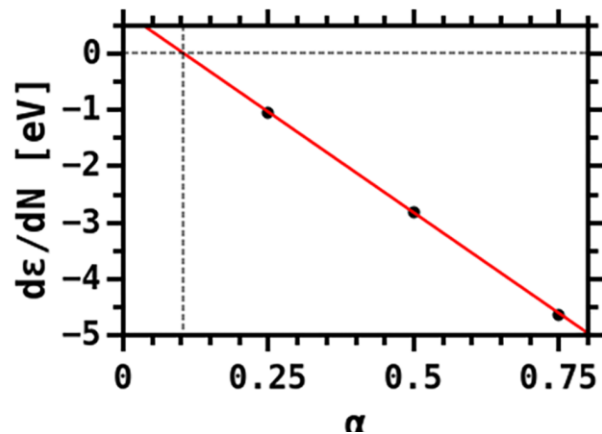
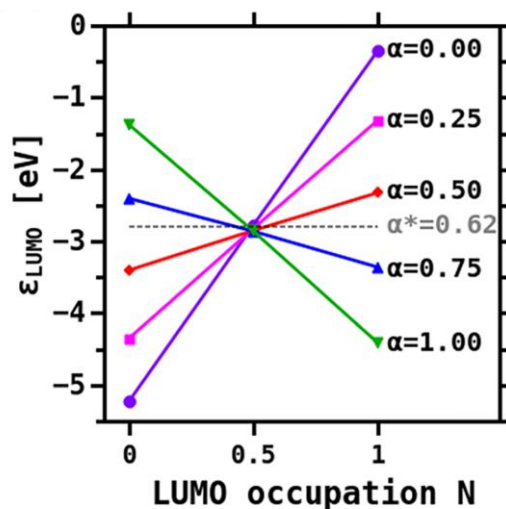
- ❖ Consider two empty states localized on two different molecules
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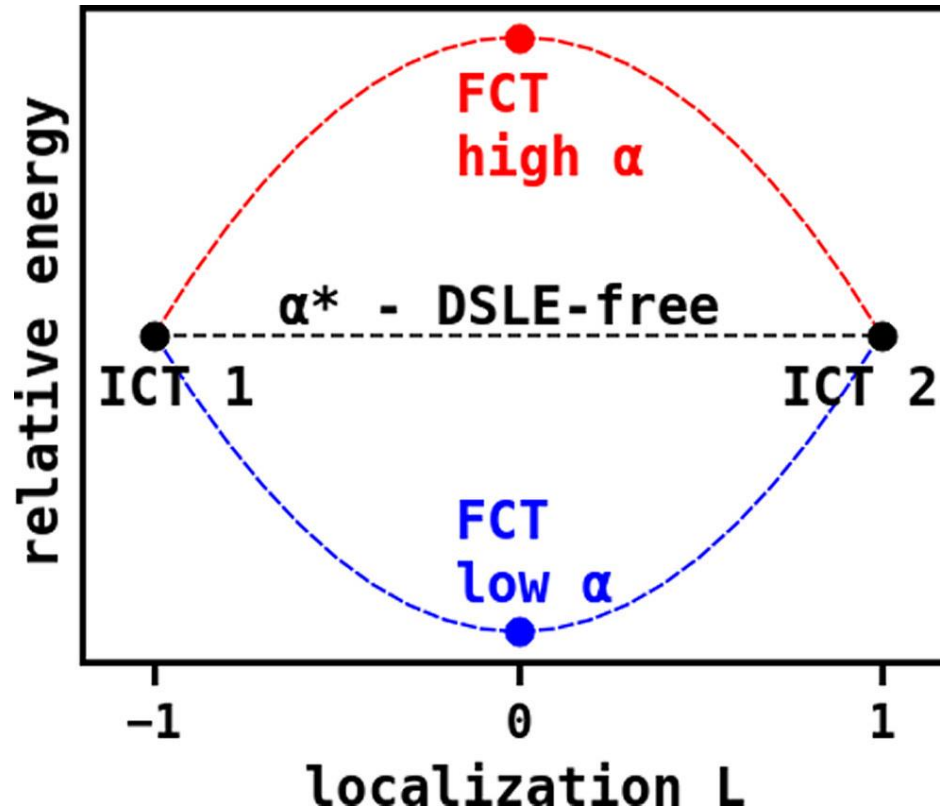


$$\frac{\delta\epsilon}{\delta n} < 0 \rightarrow \text{self-attraction}$$





- Ionization not possible \rightarrow Workarounds
- Optimal mixing depends on substrate (ϵ_r)
- Typically very small ($<25\%$)



- Small α : delocalized minimum
- Large α : localized minima, delocalized transition
- Optimal mixing: **Undecided**

The performance of semilocal DFT for CT:

➤ Overestimate CT when it is small

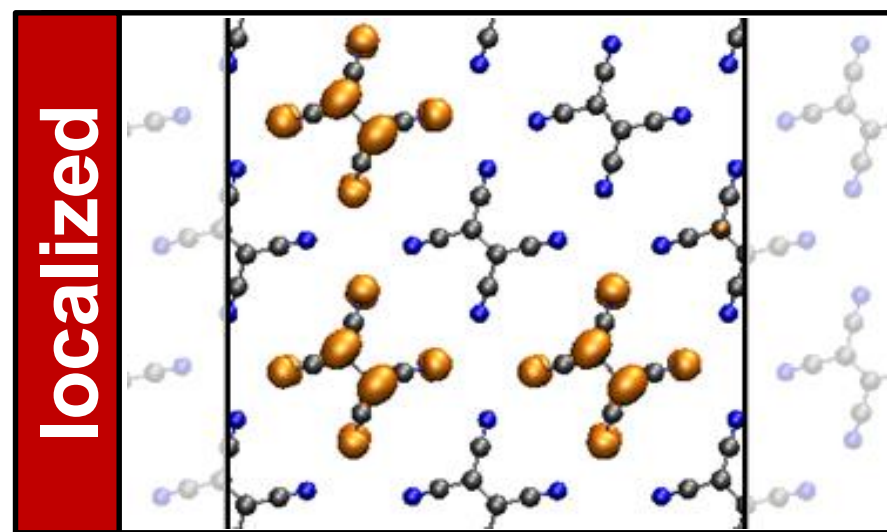
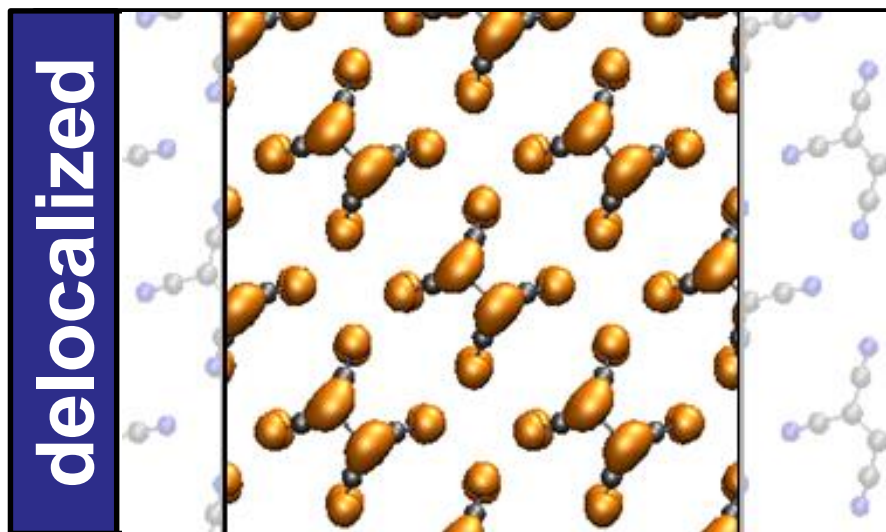
Atalla et al., Phys. Rev. B 94, 035140

➤ Underestimate CT when it is large

Hofmann et al., NJP, 15 (12), 123028, 2013

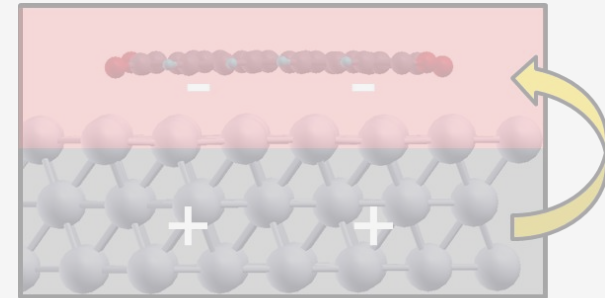
➤ Overdelocalize charge

Hofmann et al., NJP, 15 (12), 123028, 2013

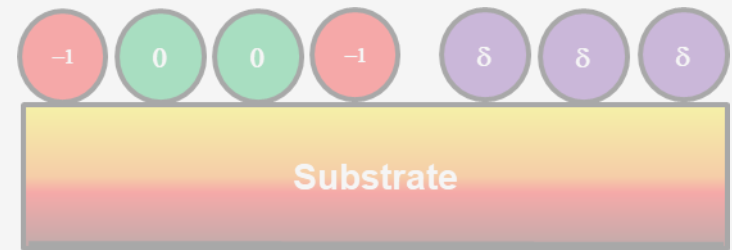


Outline: Issues with Charge Transfer

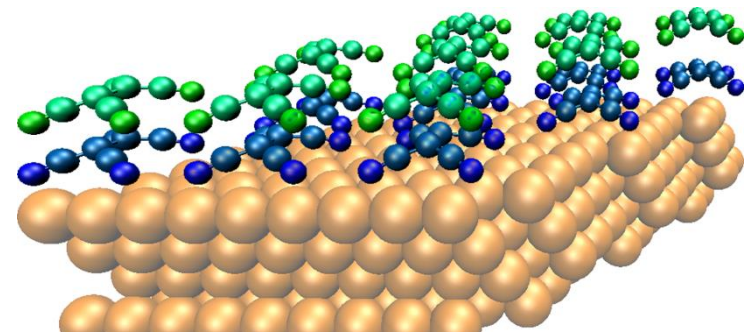
- Amount of charge transfer



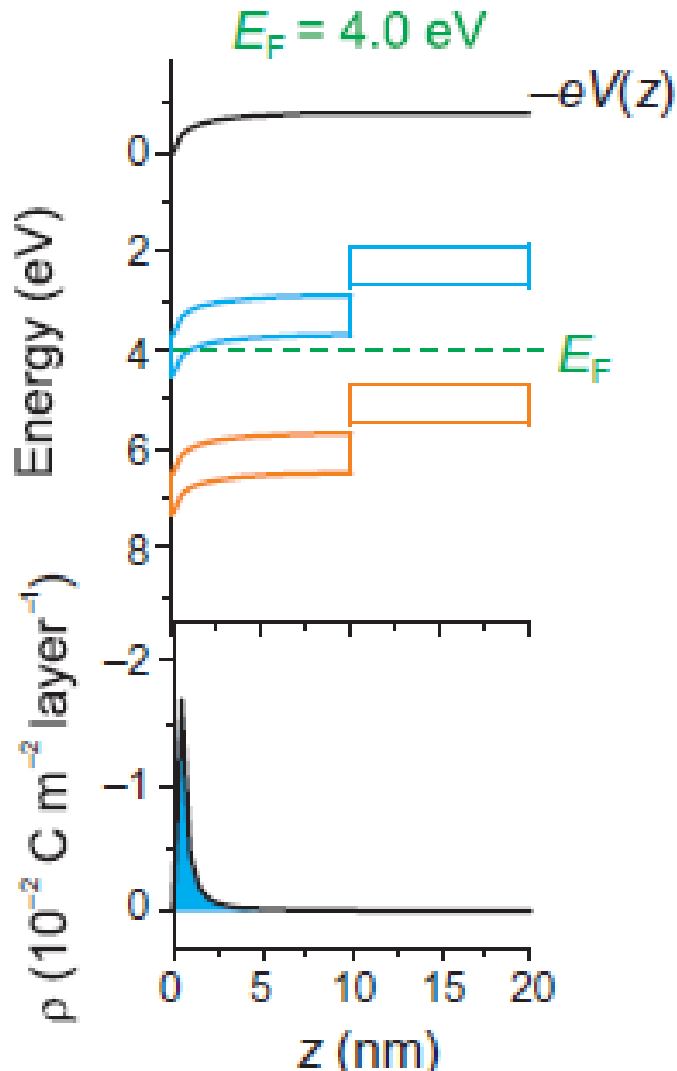
- Charge localization



- Charge transfer beyond the first layer?



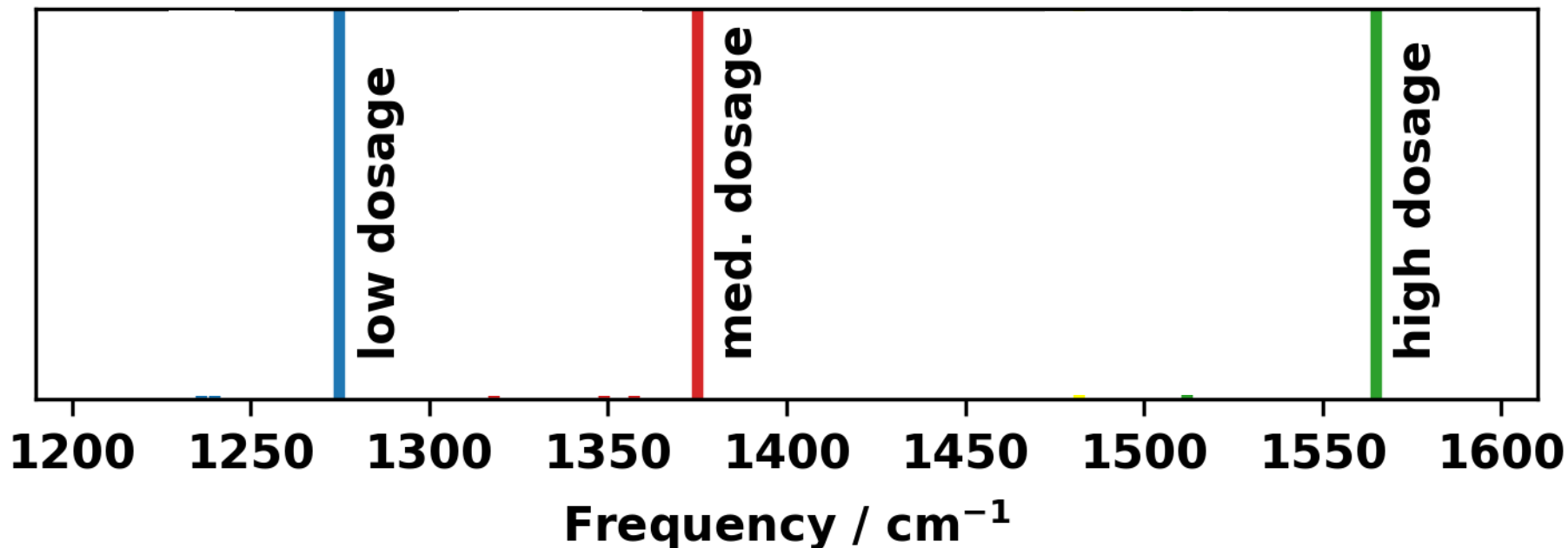
Long-ranged Charge Transfer



- Predicted by „normal“ semiconductor picture (Fermi-Dirac+Electrostatic)
- **Conflicting experiments**
 - Potential (UPS/KP)
 - Charge (Optics/Vibrations)

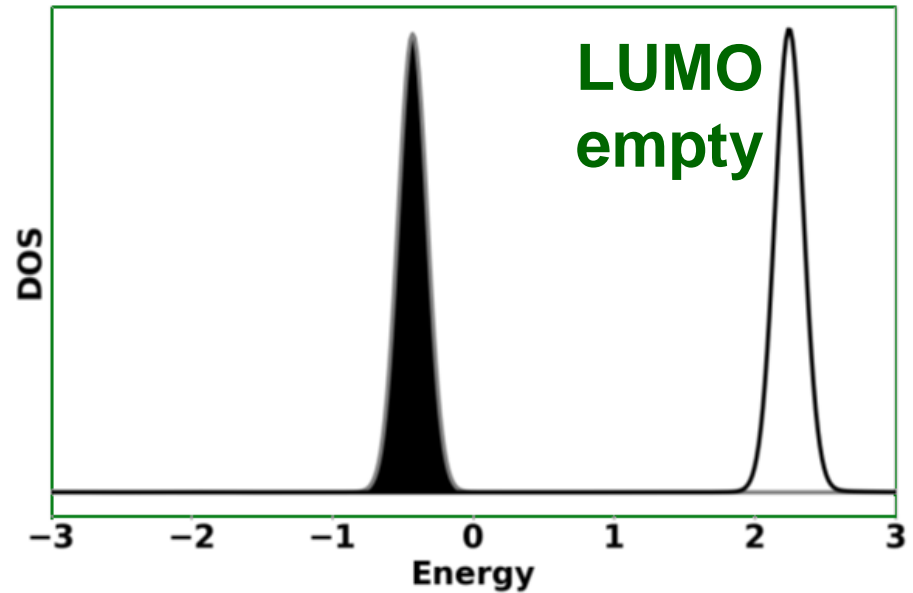
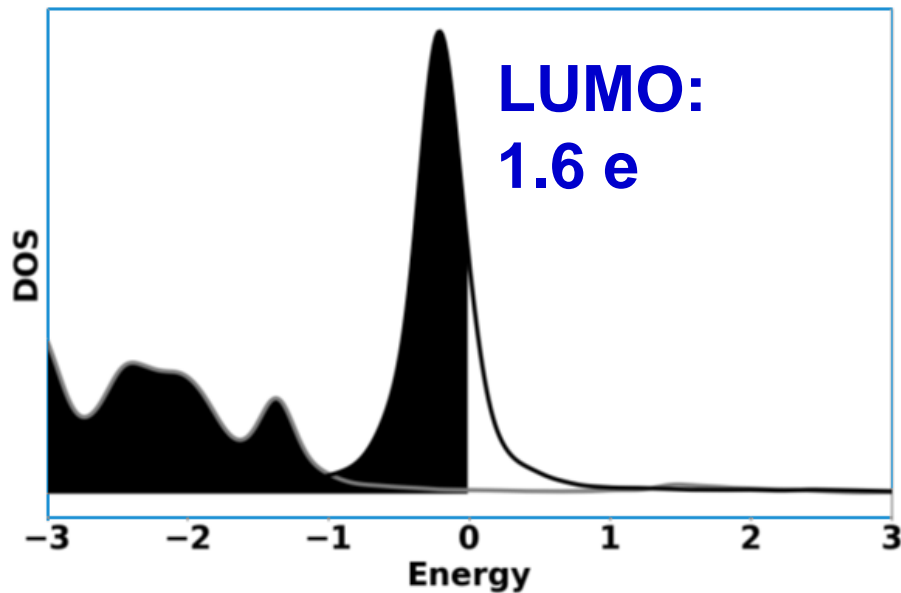
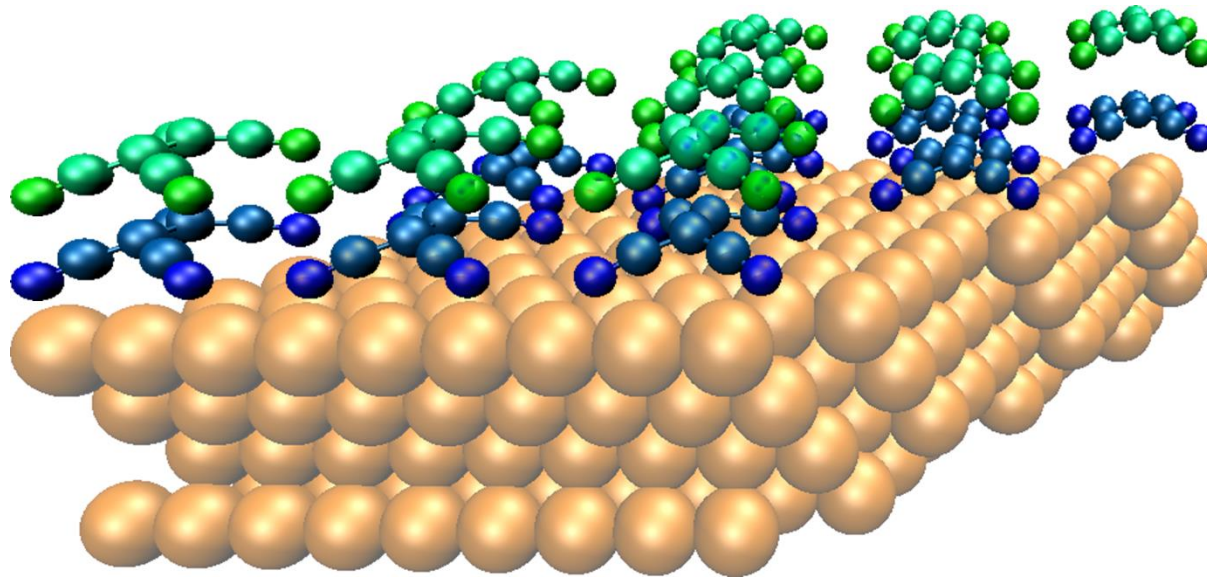
**What can we learn
from (first principles)
theory?**

A convincing experiment? TCNE/Cu(111)

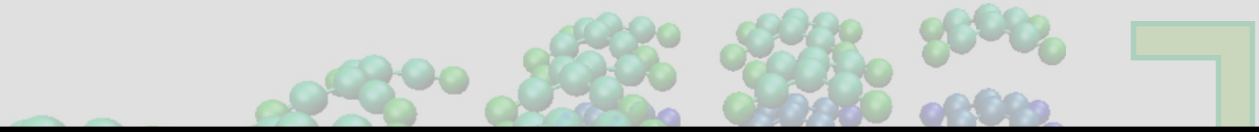


Coverage	Charge	Interpretation
Low	$>>1$	Monolayer
Medium	≈ 1	Second layer
High	0	Bulk

DFT results for ad-hoc guess structure

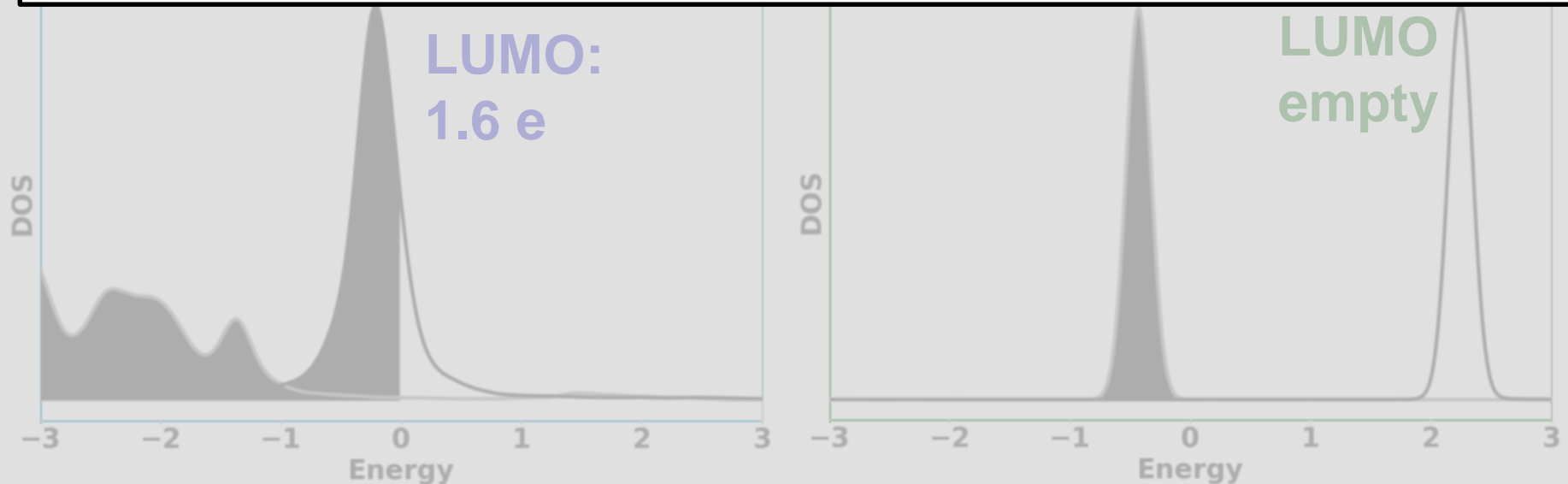


DFT results for ad-hoc guess structure



No indication for long-ranged charge transfer

No explanation of experiment



Determining bilayer structures

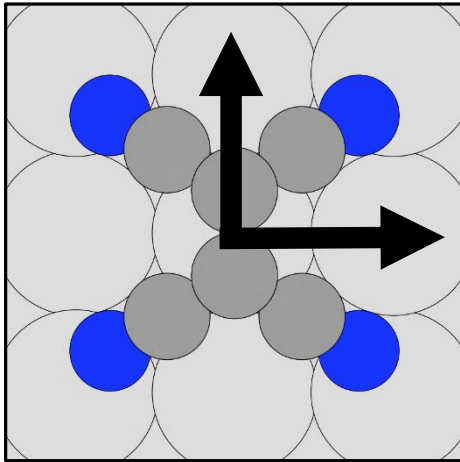
Experimentally:

- Too little material for (x-ray) diffraction
- Imaging techniques can't see „below“ bilayer
- Geometric and electronic structure rarely investigated together → comparison challenging

Computationally:

- Plausible guess structures
- Try and Error → **Confirmation Bias**
- Systematic → **Configuration Explosion**

Configurational Explosion



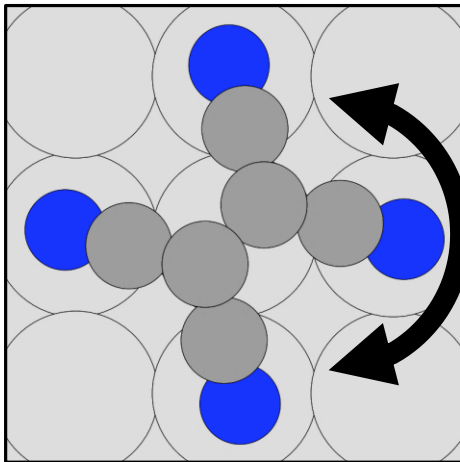
For each molecule:

Translation x: ~ 10 steps

Translation y: ~ 10 steps

Rotation: ~ 10 steps

3 mol.: $(10 \times 10 \times 10)^3 = 1 \text{ billion}$

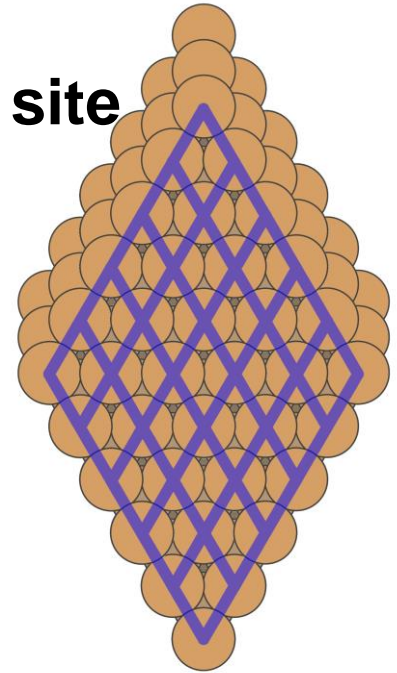
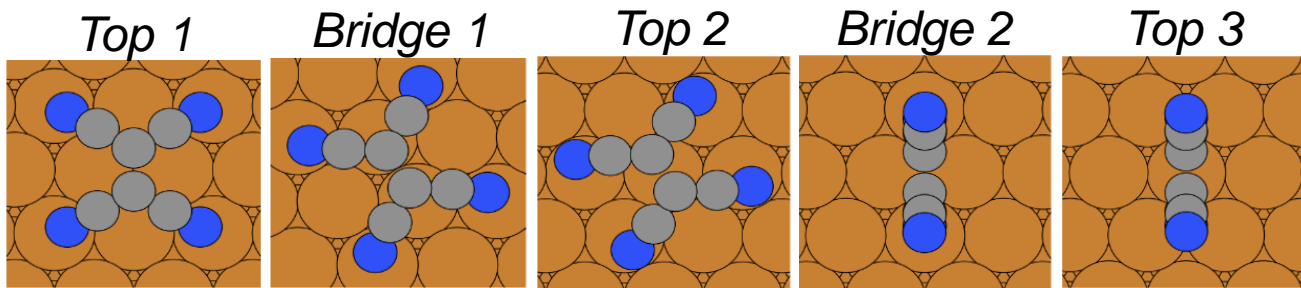


Solution: Simplify by exploiting physics at the interface

SAMPLE
=====

arXiv:1811.11702

- For commensurate interfaces, the primitive unit cell of the substrate is a natural grid.
- Each molecule sits in dedicated adsorption site



- Intermolecular interaction induces only minor changes on geometry

assemblies of
molecules that each
sit in their own
minima

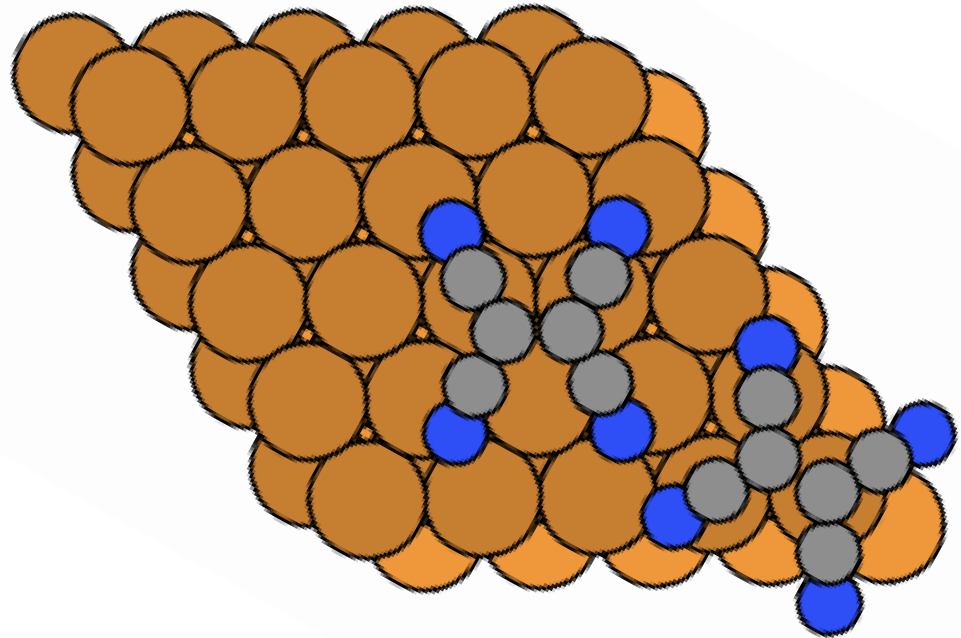


**Building block
approach**

Combine building blocks with cells

Define:

- cell size
- number of molecules
- building blocks
- min. distance



Result:

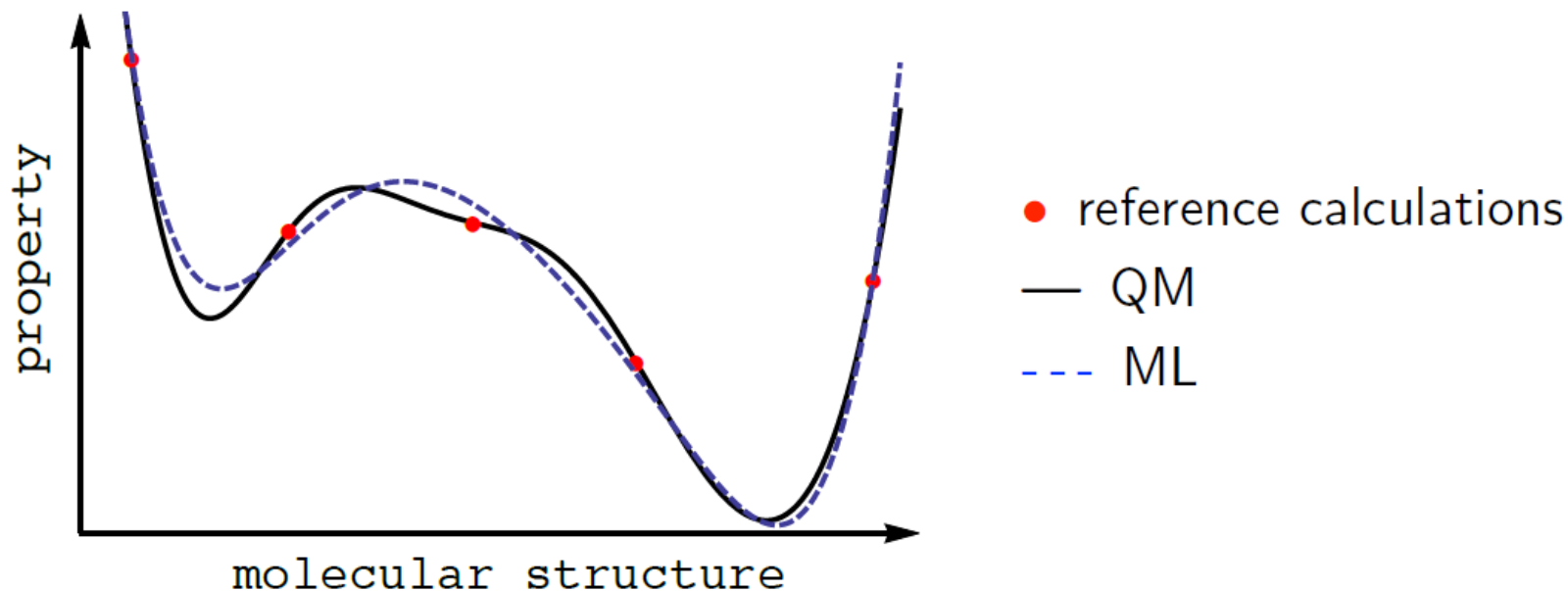
List of polymorph candidates (typically a few 100.000)

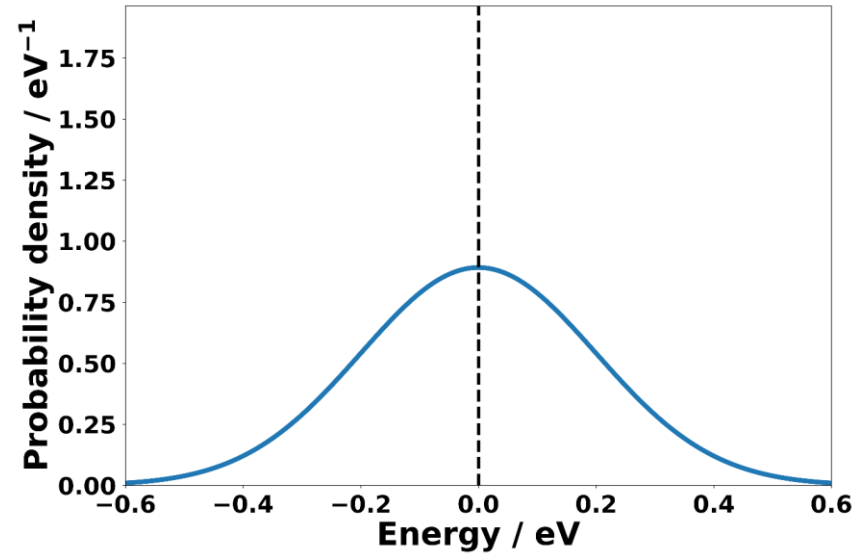
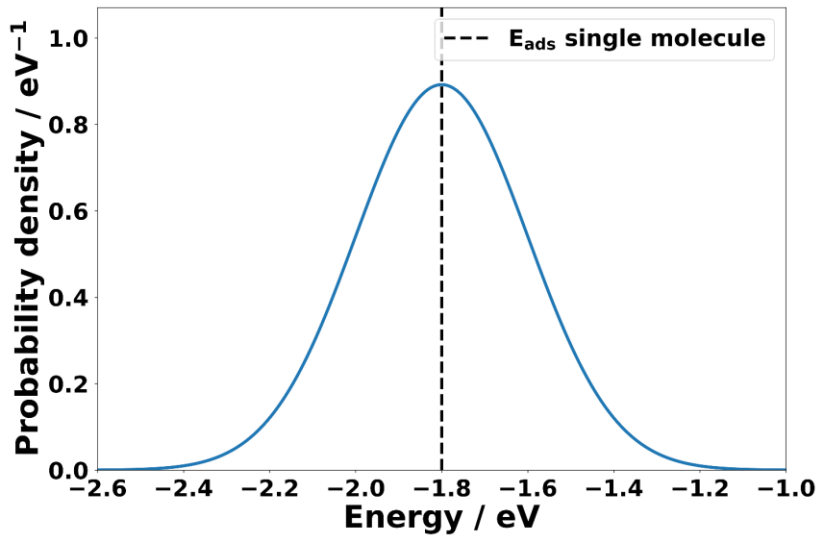
Too many for DFT  Machine Learning

The key idea of machine learning

- Exploit **redundancies** in **related** calculations
- **Interpolate** between calculations
- **Smoothness** assumption (regularization)

Smoothness assumption can be based on known physics at this kind of interfaces

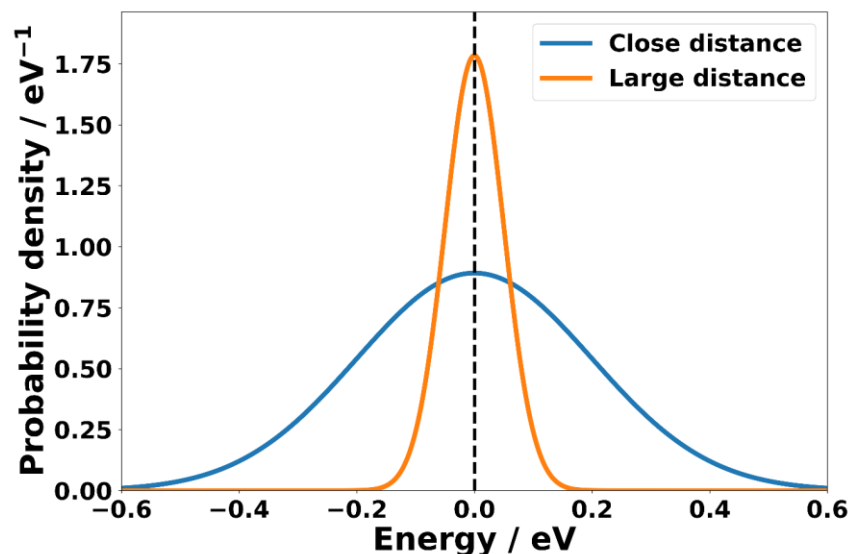
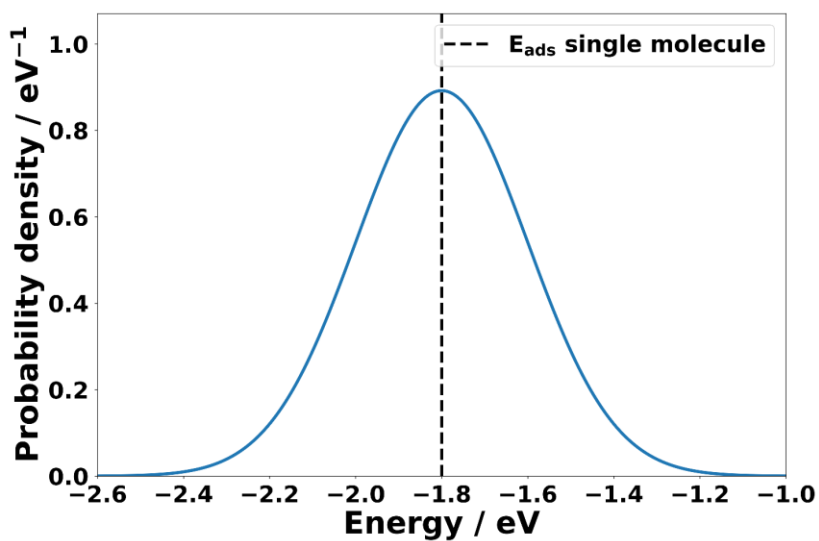




$$\mathbf{E} = \sum_{\mathbf{i}} n_{\mathbf{i}} U_{\mathbf{i}} + \sum_{\mathbf{p}} n_{\mathbf{p}} V_{\mathbf{p}}$$

Similar to isolated molecule

Individual terms are small



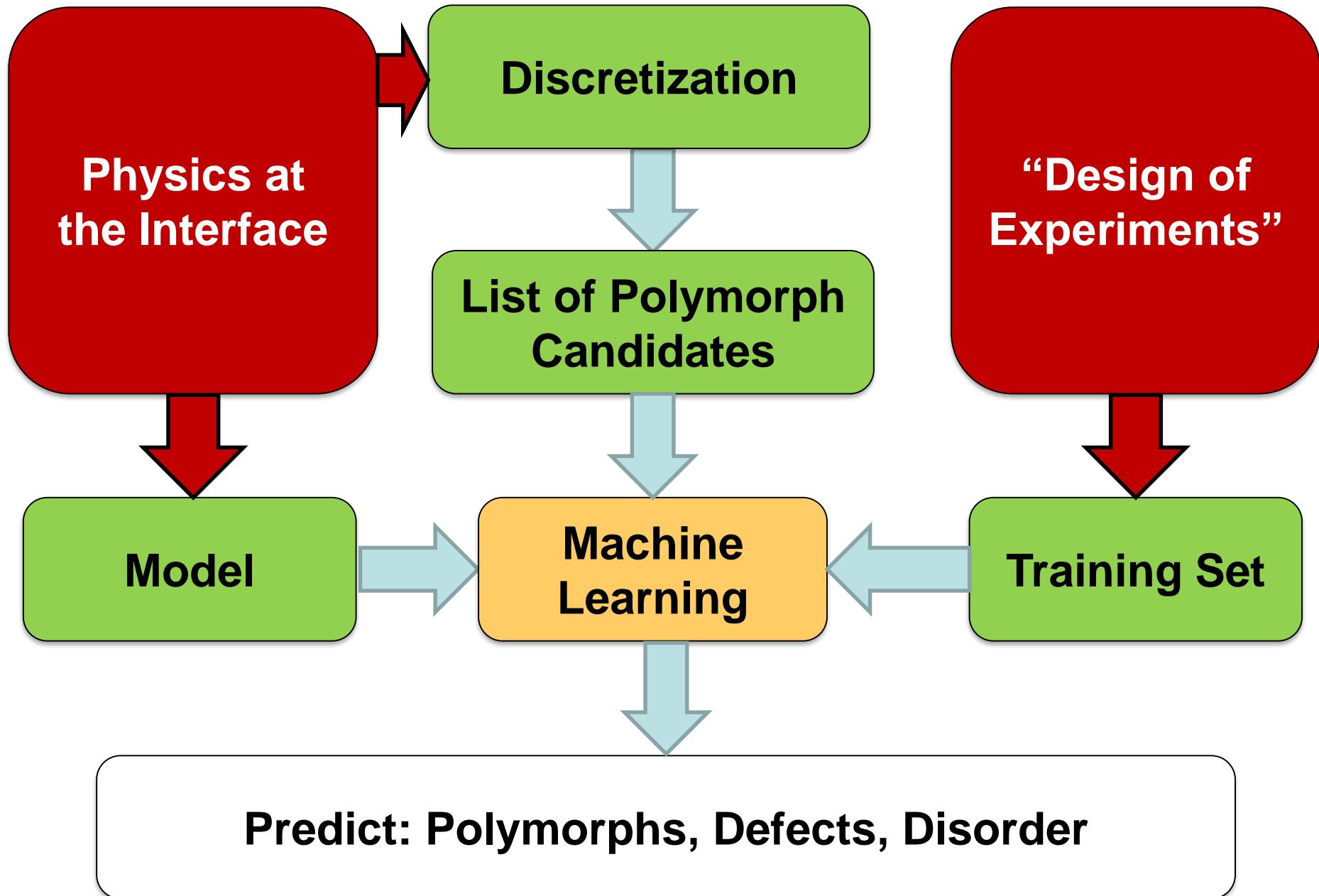
$$\mathbf{E} = \sum_{\mathbf{i}} n_{\mathbf{i}} \mathbf{U}_{\mathbf{i}} + \sum_{\mathbf{p}} n_{\mathbf{p}} \mathbf{V}_{\mathbf{p}}$$

Similar to isolated molecule

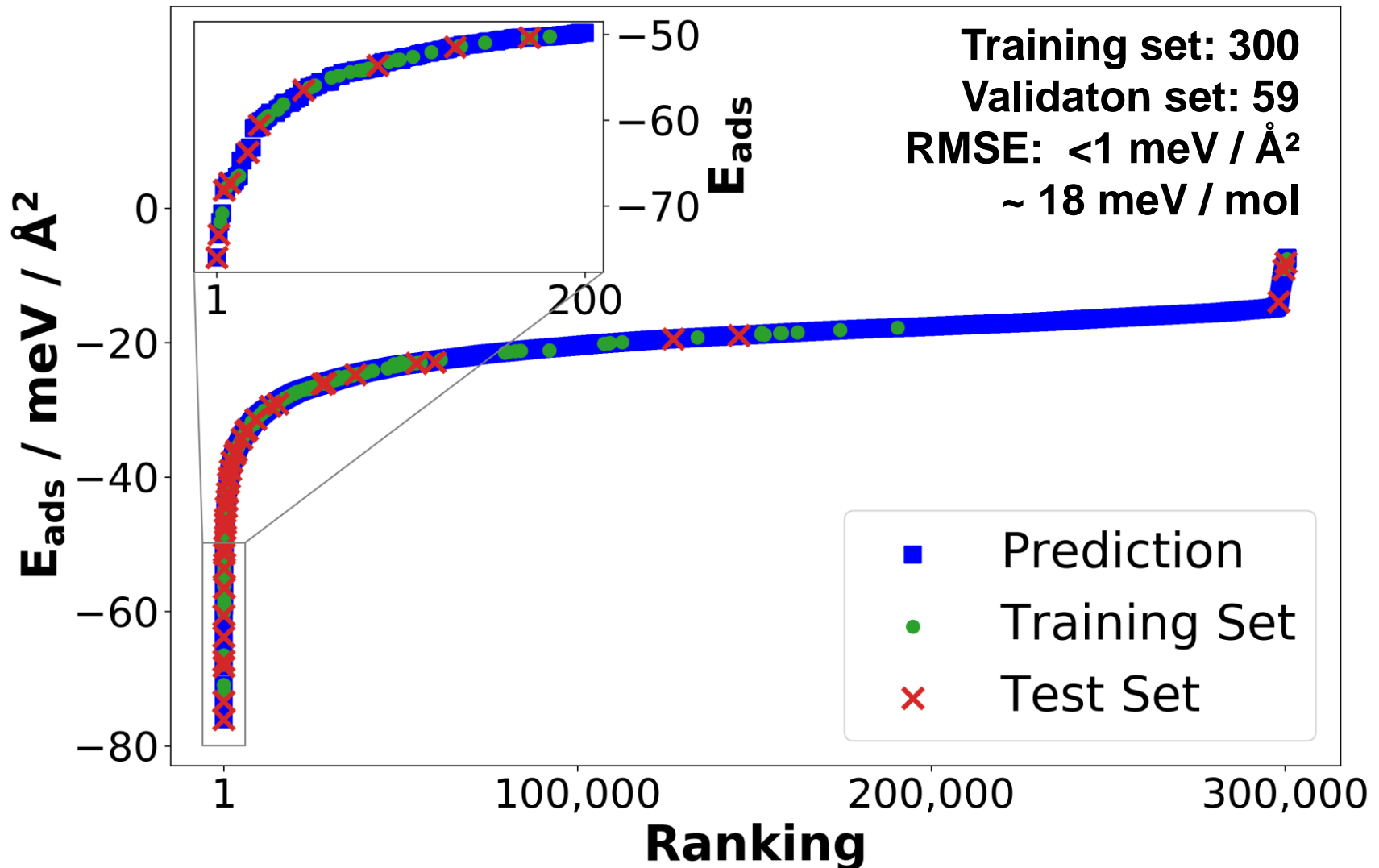
Individual terms are small

Larger distance → less interaction

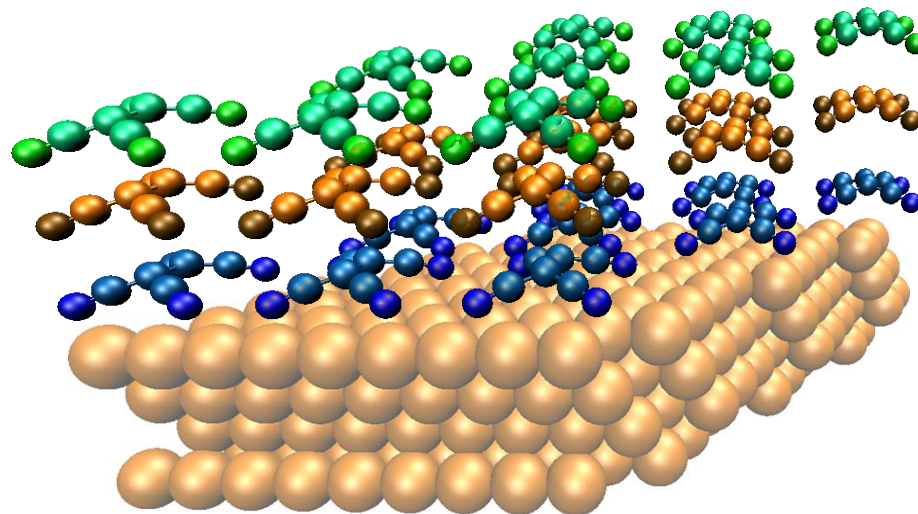
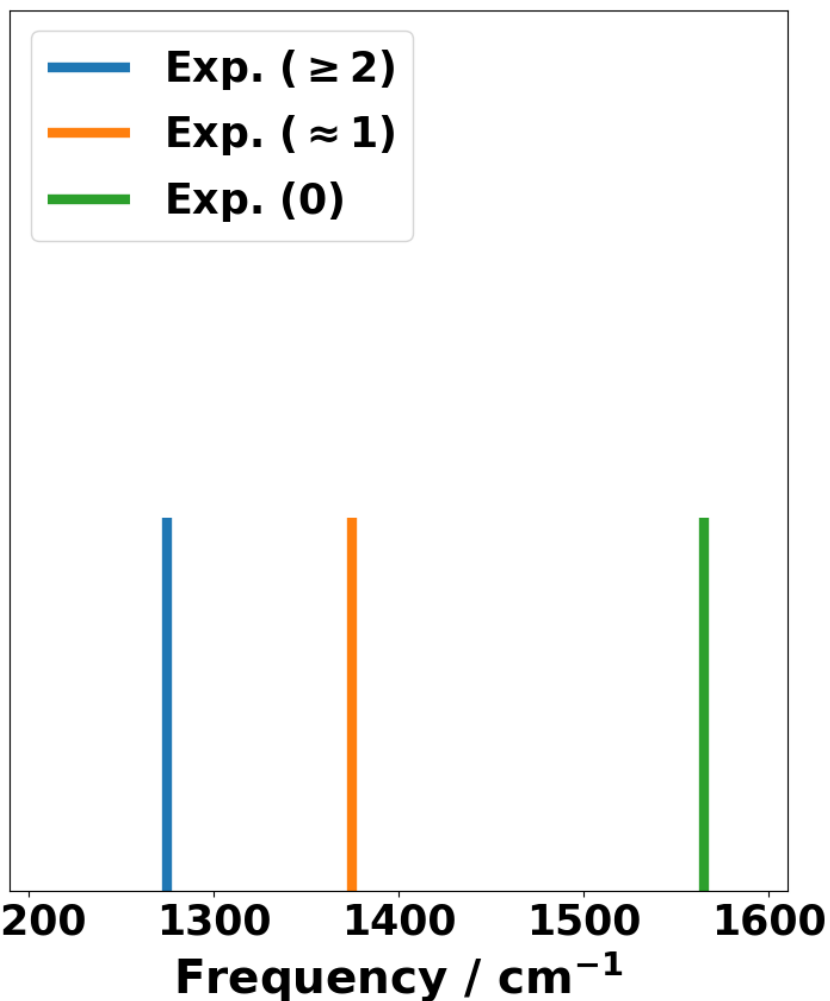
Similar structures → similar interaction



Machine Learning Performance



Back to the original question

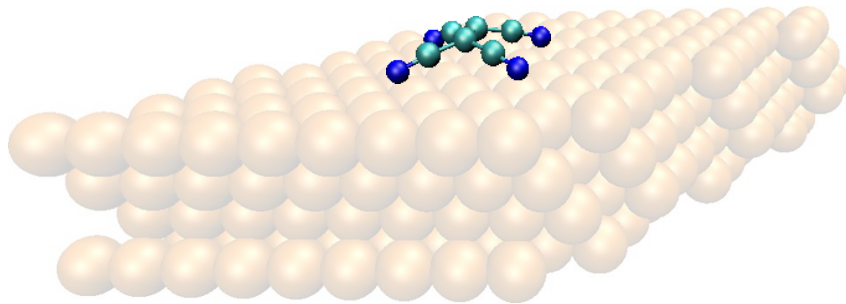


Where is the singly charged species?

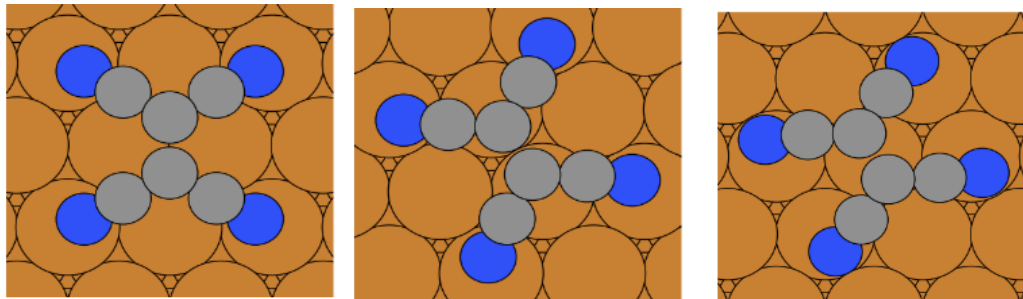
What is the correct structure?

Building blocks

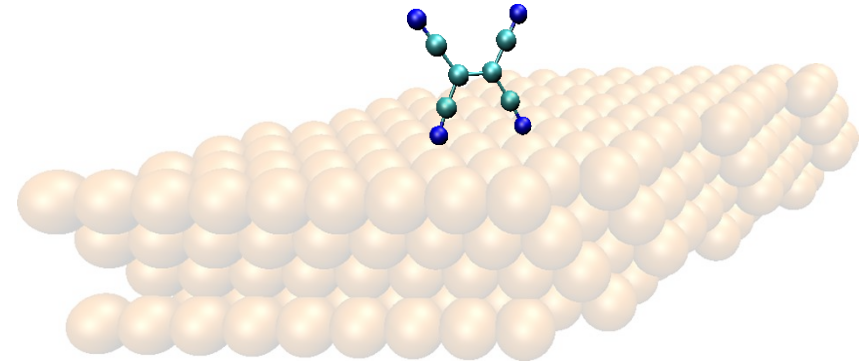
Flat lying



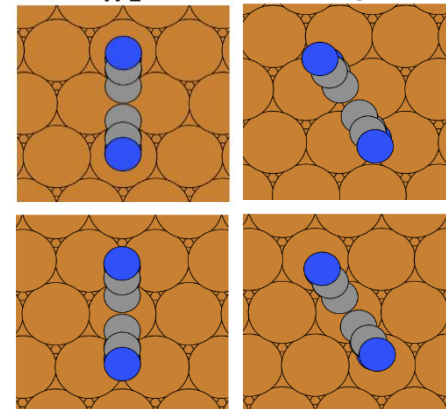
Energetically more favorable



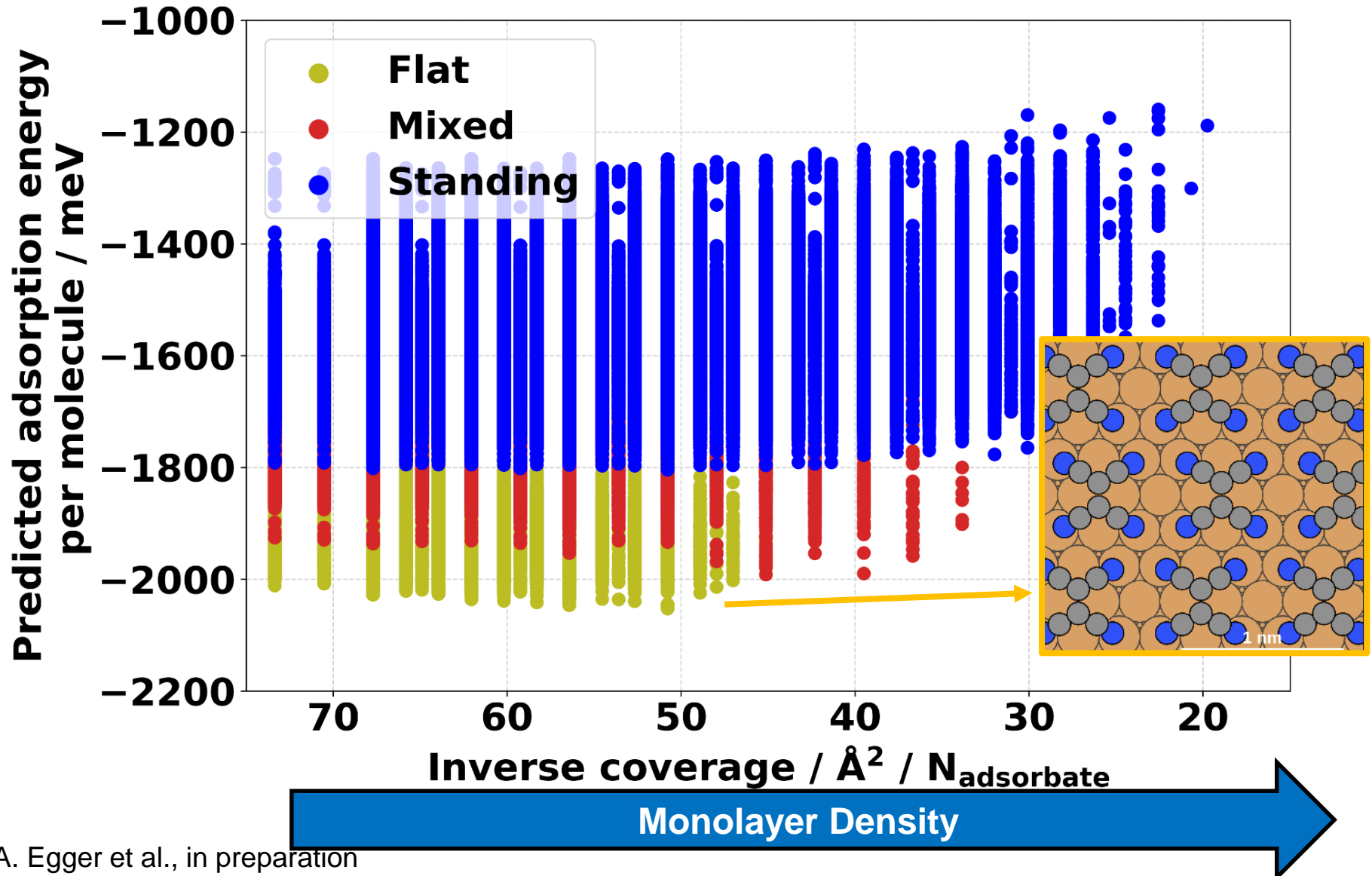
Standing

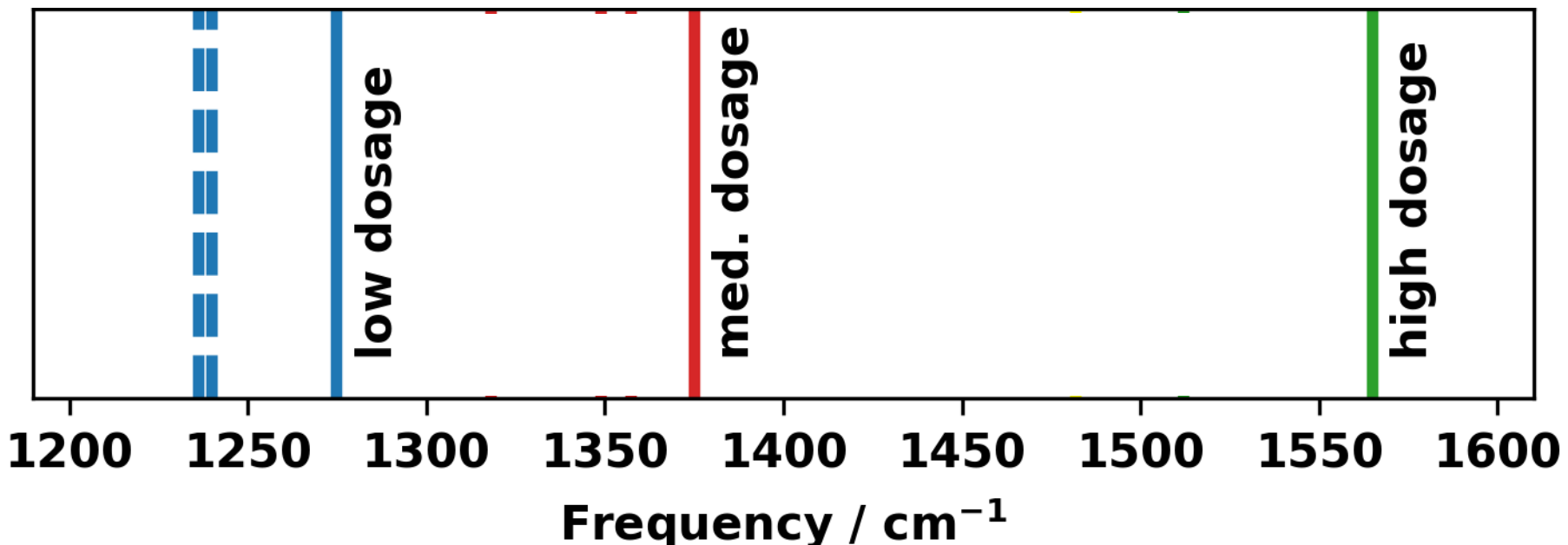
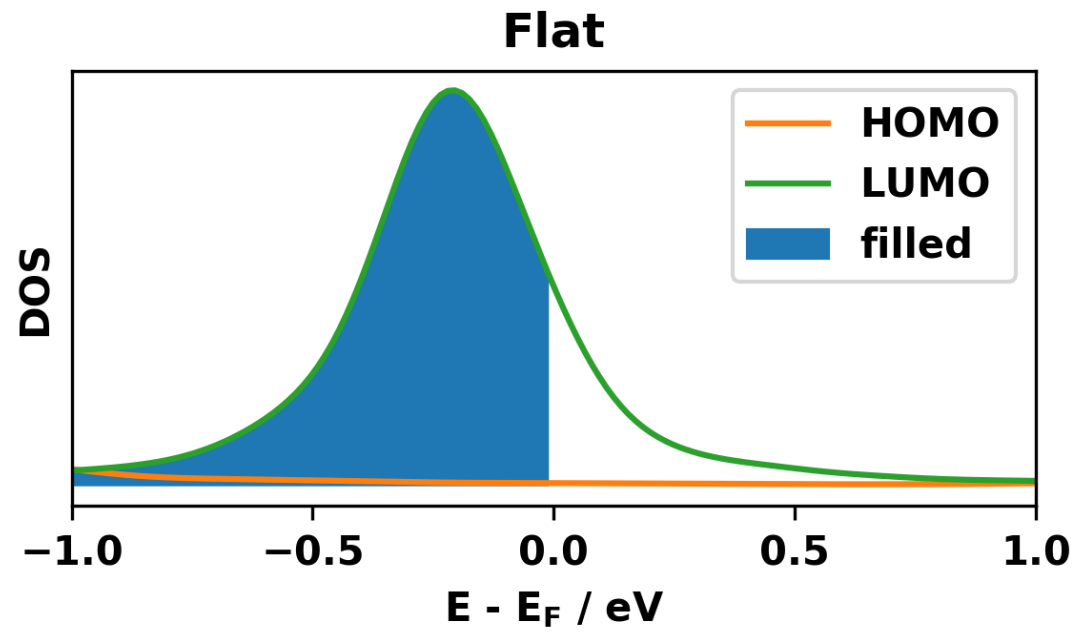
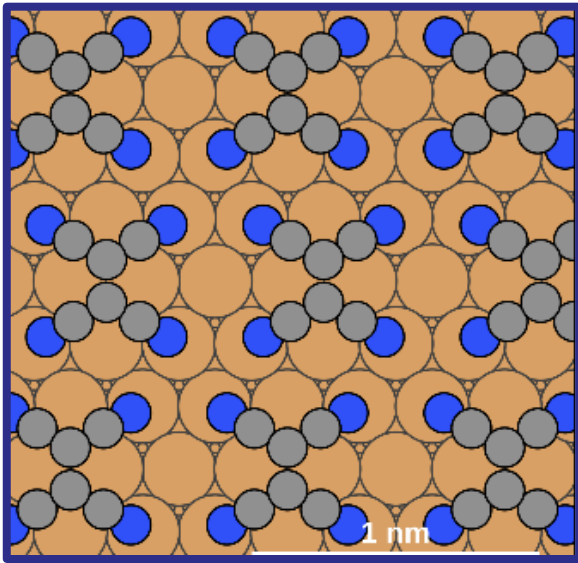


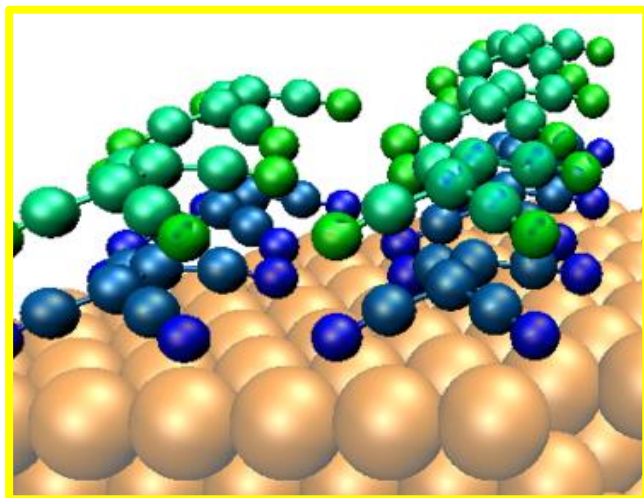
Smaller footprint



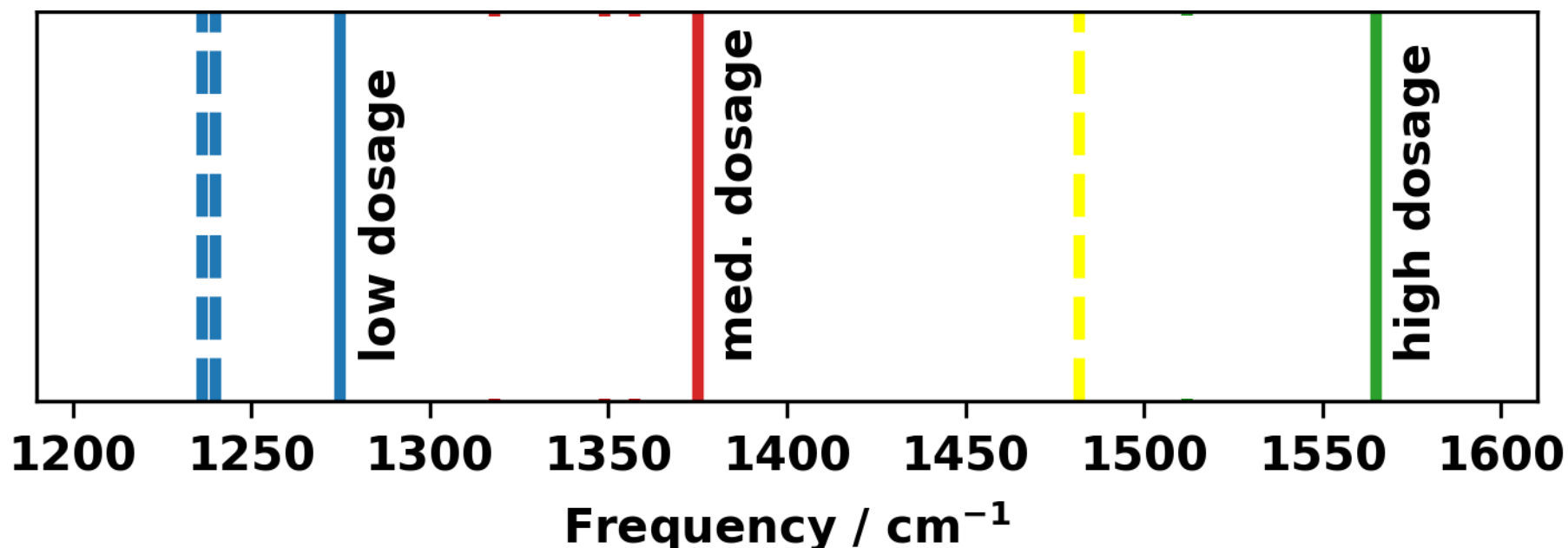
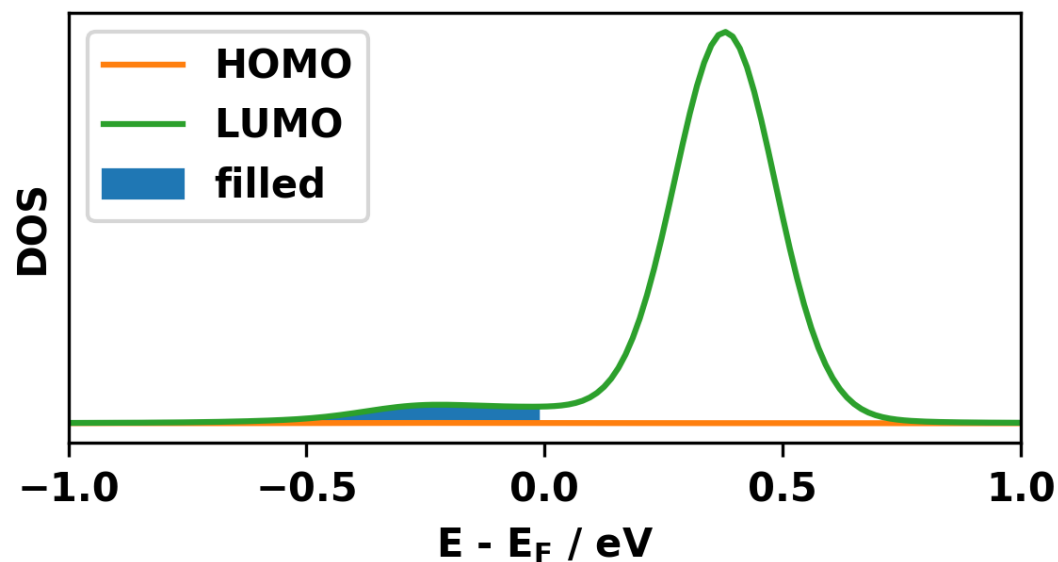
Submonolayer: best energy per molecule



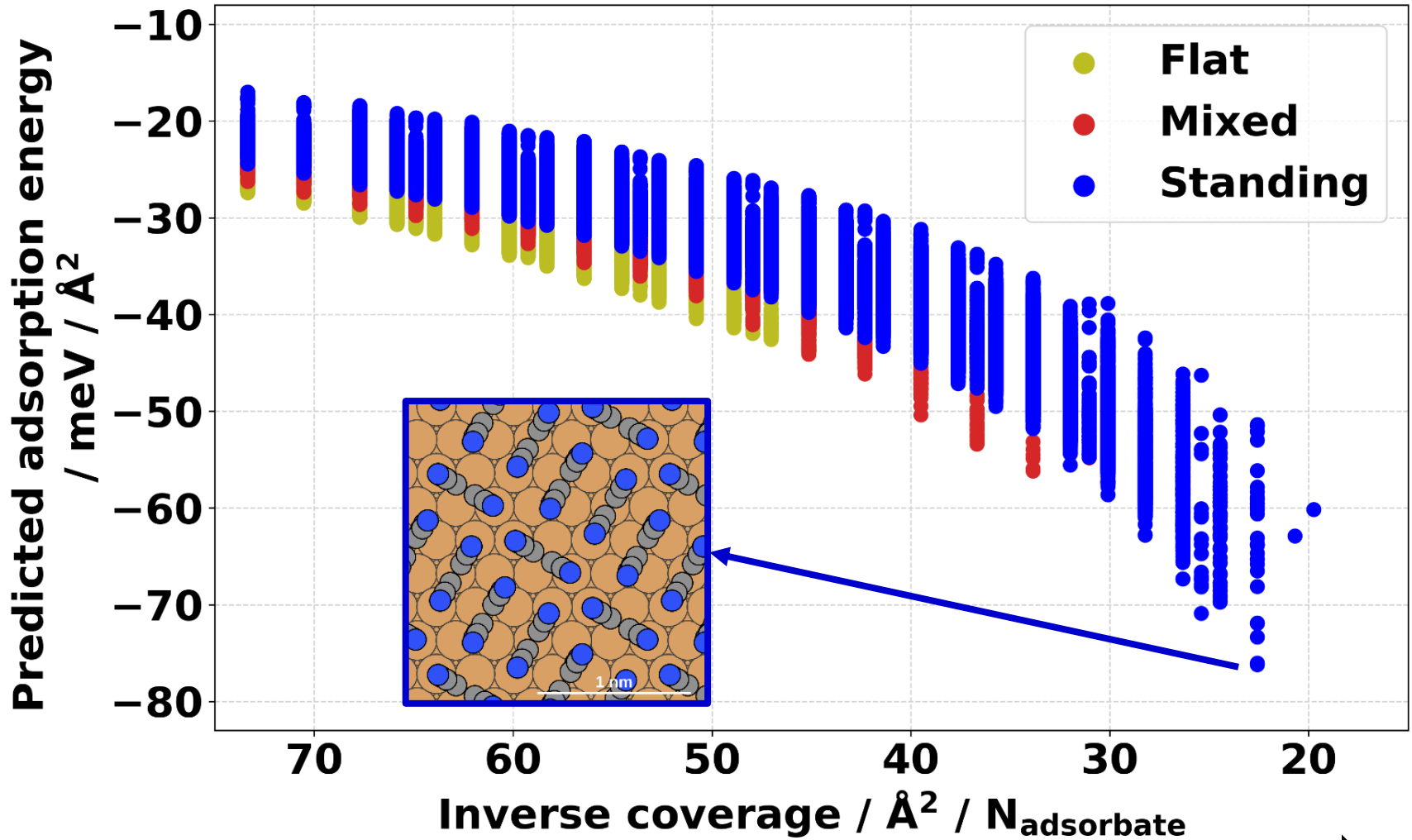


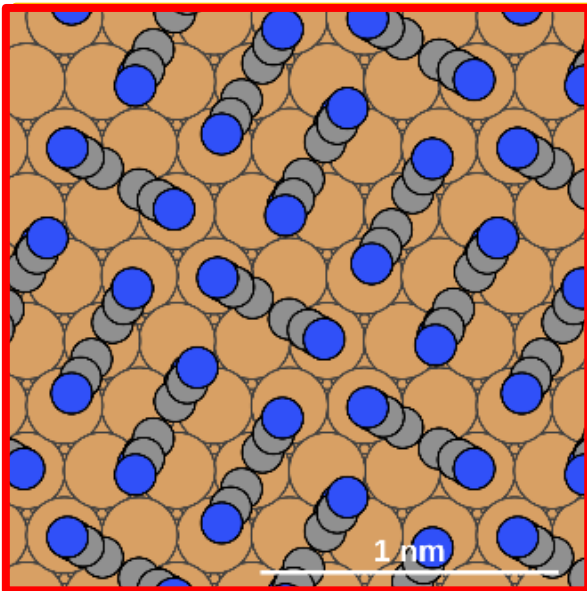


Flat-on-Flat

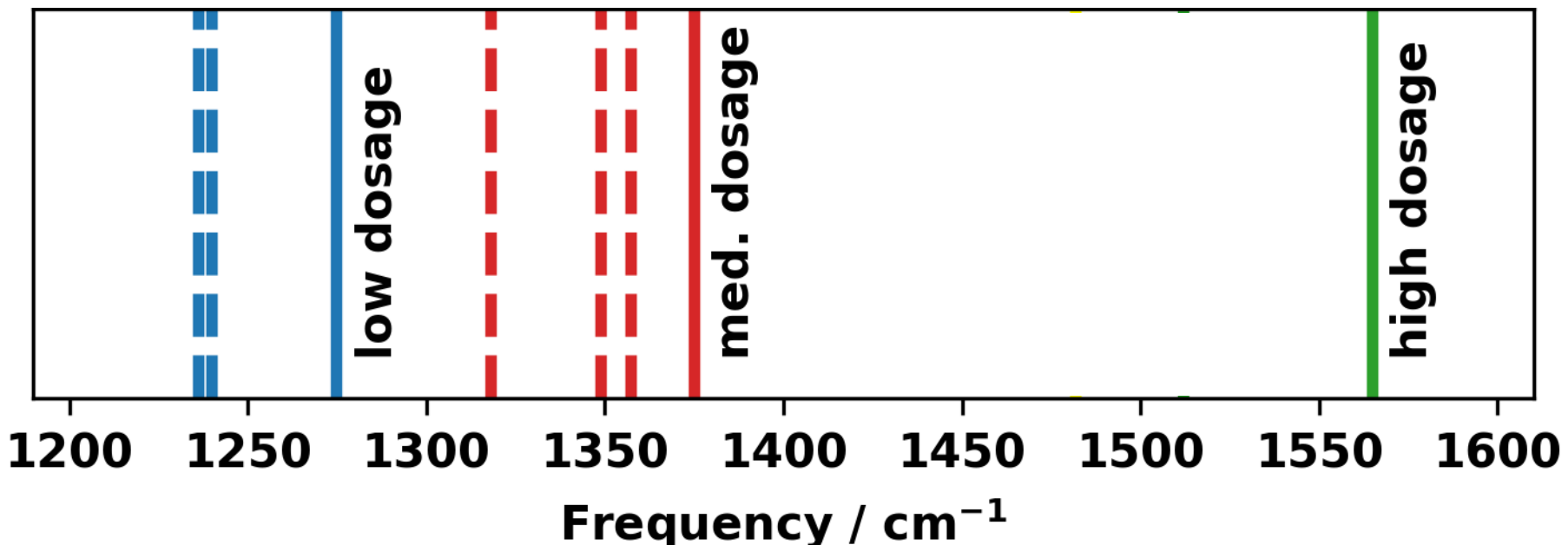
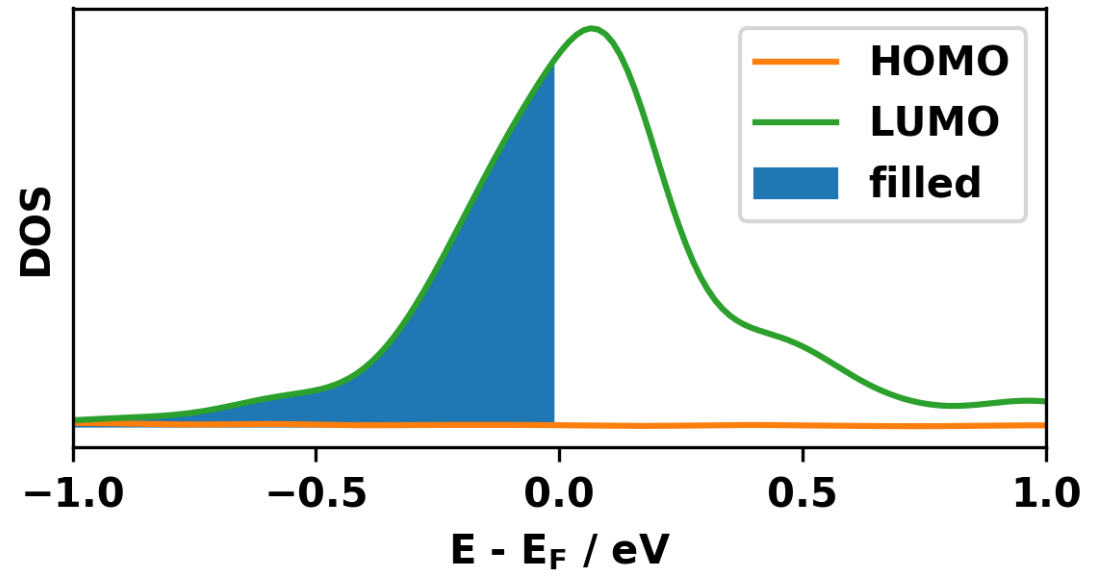


Closed layer: best energy per area

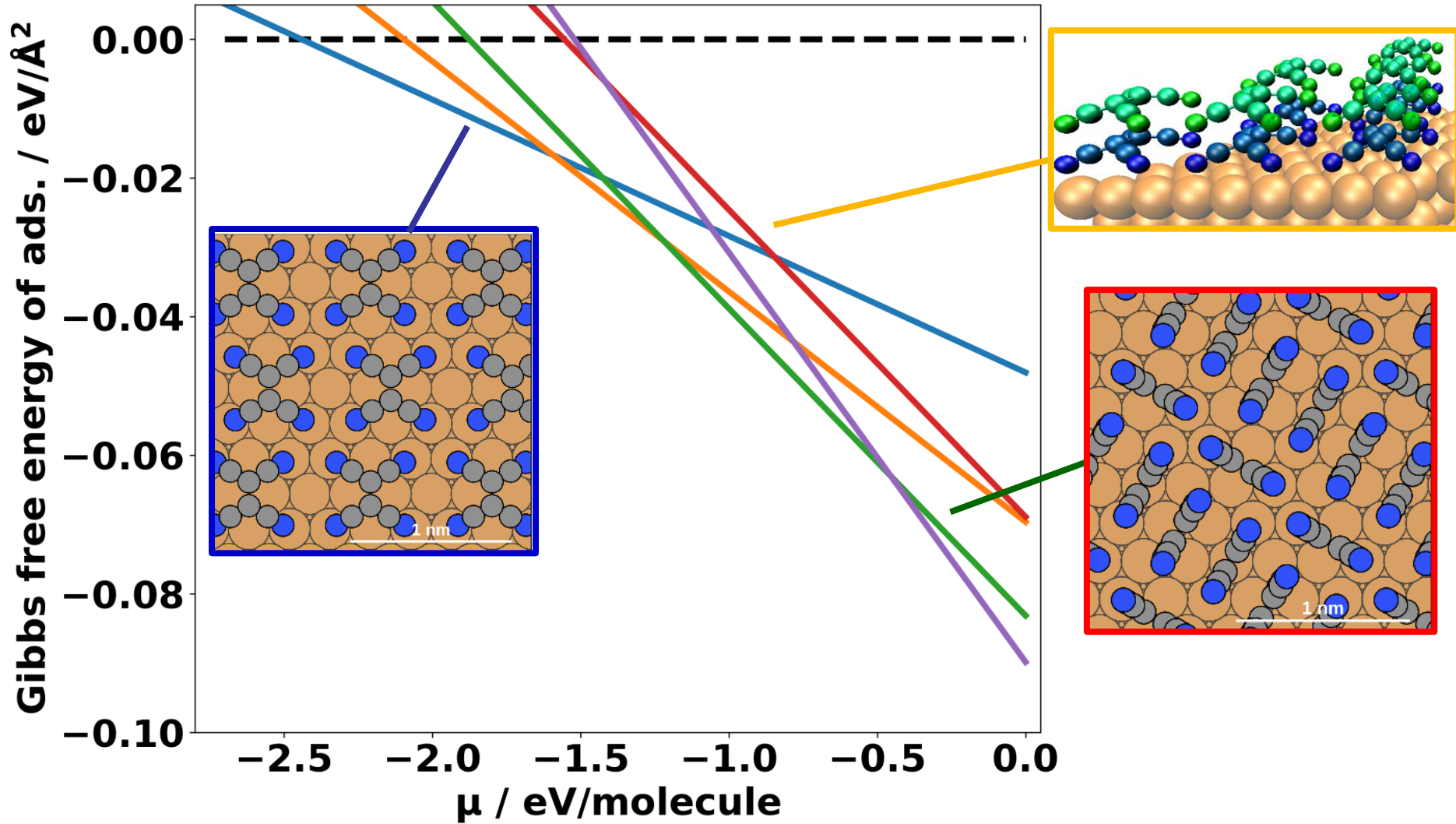




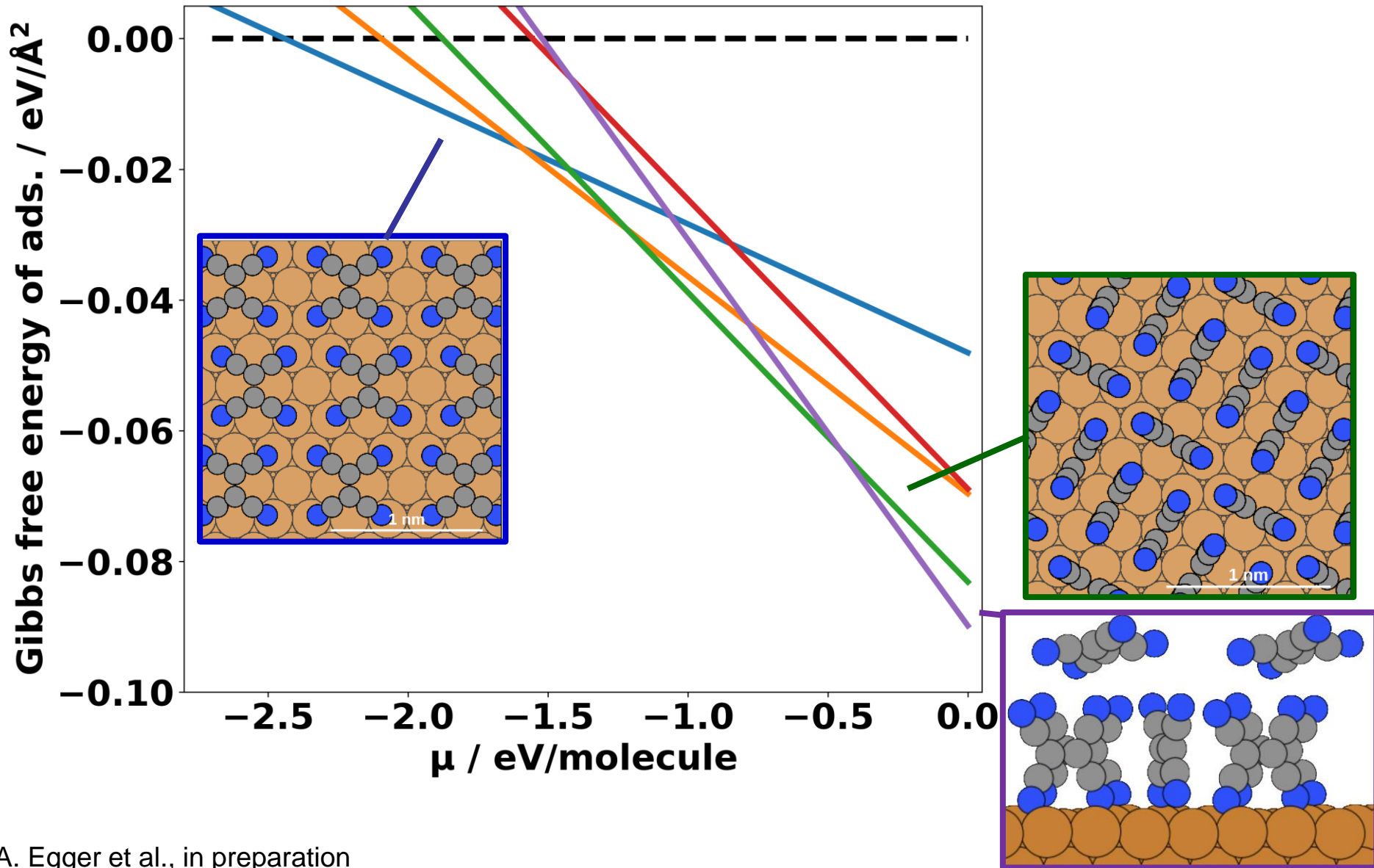
Standing

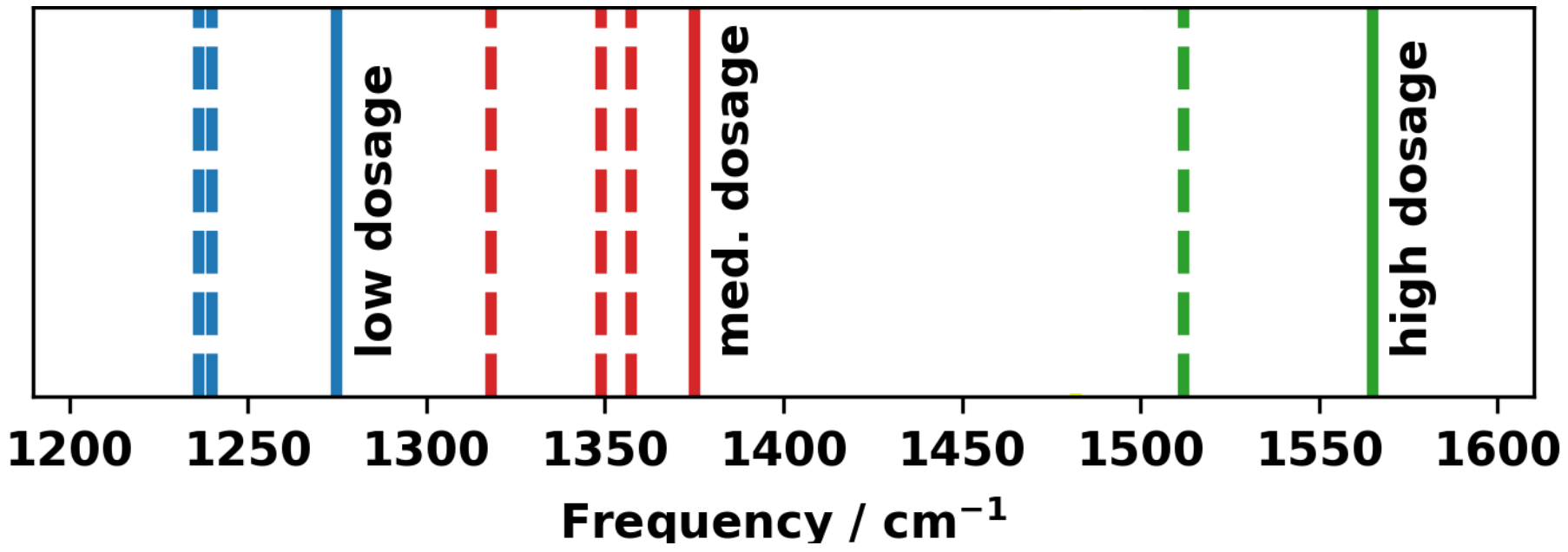
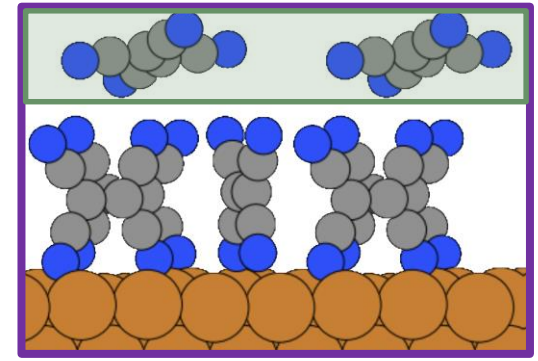
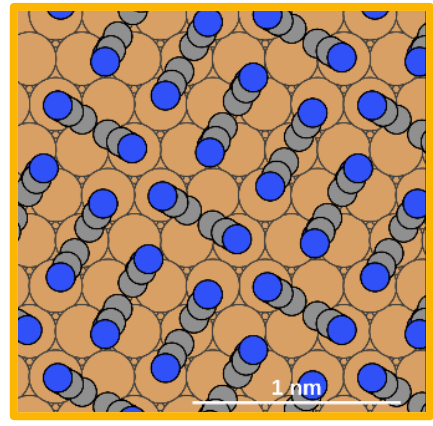
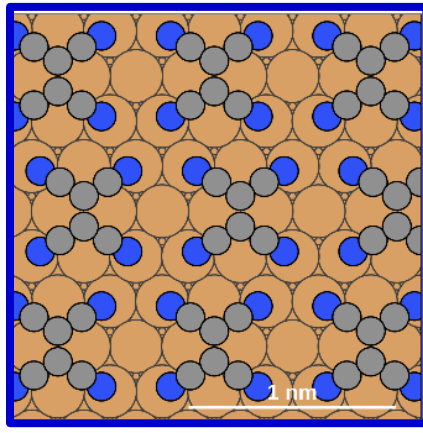


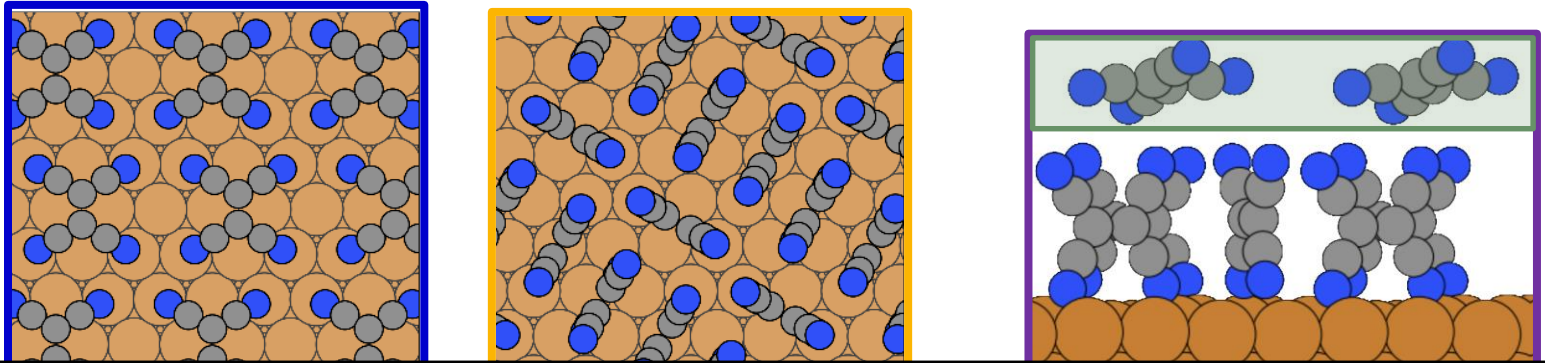
Ab-initio thermodynamics



Ab-initio thermodynamics



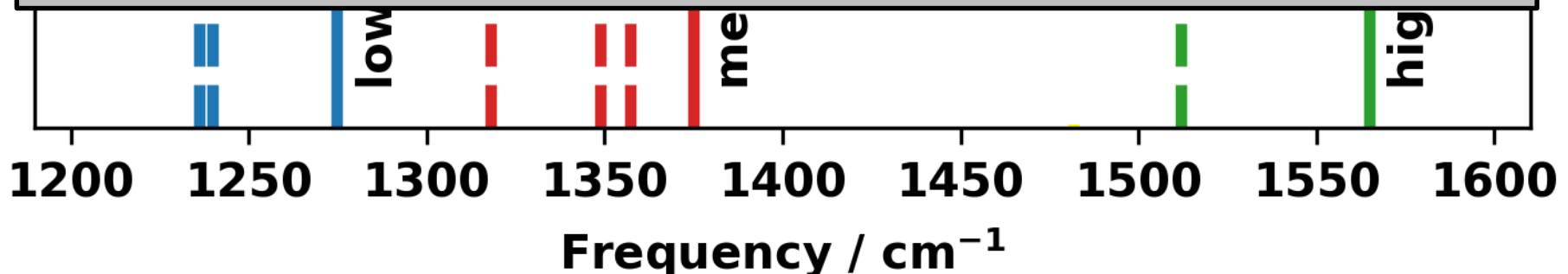




Conclusion:

There is a phase transition from lying to standing before the formation of the second layer

„Standing“ obtains less charge than „Lying“



Summary

DFT indispensable for materials science

- Prediction of new materials
- Understanding relevant processes

Hard Problems \neq unsolvable

No functional is perfect

- Understand limitations and accuracy
- (Numerical) Reasons for those