

The Numerical Assembly Technique for arbitrary planar beam systems based on an improved homogeneous solution

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- 1 Overview
- 2 Arbitrary planar beam systems
- 3 Improved homogeneous solution
- 4 Numerical Assembly Technique
- 5 Results

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Content

- General formulation of the boundary and interface conditions for arbitrary planar beam structures
- Study of the homogeneous solution of the Timoshenko beam
- Generating the system matrix by means of the Numerical Assembly Technique
- Determining the eigenvalues and eigenvectors of the system matrix

Aim and scope

- Natural Frequencies of any continuous planar beam structure considering the Timoshenko beam theory
- Analytical solutions of the differential equation*s
- Improvement of the computational performance by use of an improved homogeneous solution

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1 Overview

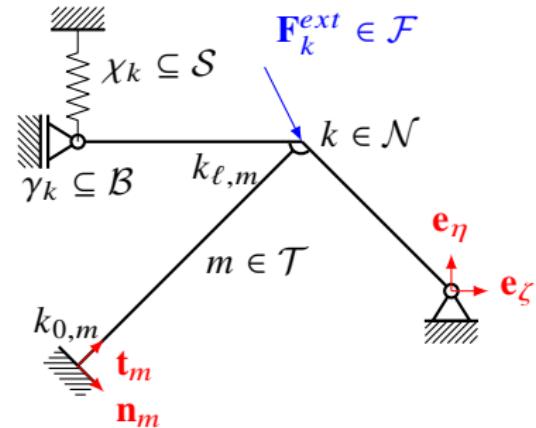
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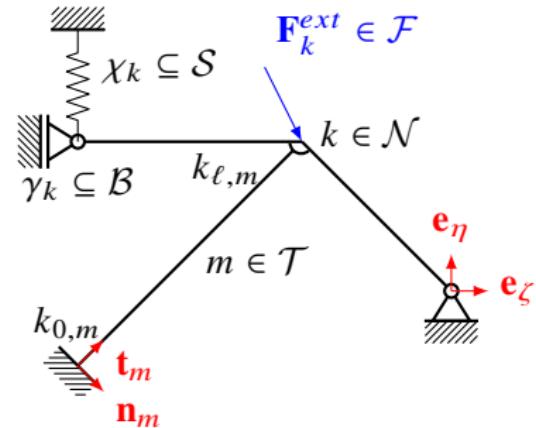
- \mathcal{N} ... set of nodes
- \mathcal{S} ... set of springs
- \mathcal{T} ... set of beams
- \mathcal{F} ... set of external loads
- \mathcal{B} ... set of bearings



- Each beam $m \in \mathcal{T}$ has
 - two nodes $k_{0,m}, k_{\ell,m} \in \mathcal{N}$
 - parameters ρ_m, E_m, I_m, A_m
 - rigid or hinged at k_0, k_ℓ
 - $\mathbf{t}_m = \frac{k_{\ell,m} - k_{0,m}}{\|k_{\ell,m} - k_{0,m}\|}$
 - $\mathbf{n}_m \perp \mathbf{t}_m$

- Each node $k \in \mathcal{N}$ has
 - set of beams $\beta_k \subseteq \mathcal{T}, |\beta_k| \geq 1$
 - set of springs $\chi_k \subseteq \mathcal{S}, |\chi_k| \in \{0, 3\}$
 - external loads $\mathbf{F}_k^{ext} \in \mathcal{F}$
 - external moments $M_k^{ext} \in \mathcal{F}$
 - bearing $\gamma_k \subseteq \mathcal{B}, |\gamma_k| \in \{0, 1\}, \mathbf{e}_\zeta \perp \mathbf{e}_\eta$

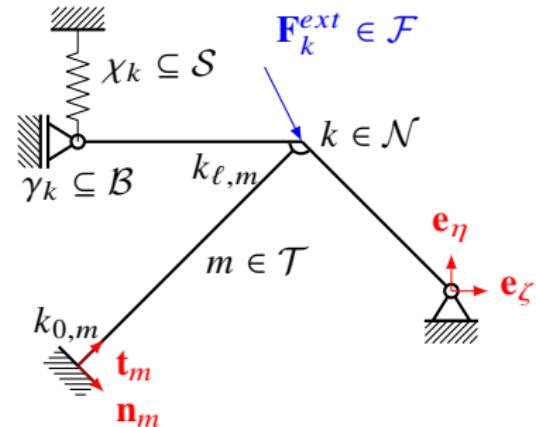
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- The rotation of the beam axis ψ_m
- The resulting deformation

$$\mathbf{U}_m(\xi) = U_m(\xi) \cdot \mathbf{t}_m + W_m(\xi) \cdot \mathbf{n}_m$$

- Internal bending moment M_m
- The resulting internal force

$$\mathbf{S}_m(\xi) = N_m(\xi) \cdot \mathbf{t}_m + V_m(\xi) \cdot \mathbf{n}_m$$

Preliminaries

- The rotation of the beam axis ψ_m
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■ Kinetics

$$\mathbf{F}_k = \sum_{m \in \beta_k} \mathbf{S}_{m|k} + \mathbf{B}_k^L + \mathbf{F}_k^{ext} + \mathbf{F}_k^L = \mathbf{0} \quad \text{free vib.} \rightarrow \mathbf{F}_k^{ext} = \mathbf{0}$$

$$M_k = \sum_{m \in \beta_k} M_{m|k} + B_k^R + M_k^{ext} + M_k^R = 0 \quad \text{free vib.} \rightarrow M_k^{ext} = 0$$

■ Springs

$$\mathbf{F}_k^L = -c_L(\mathbf{U}_k \cdot \boldsymbol{\zeta}) \cdot \boldsymbol{\zeta} \quad c_L \dots \text{longitudinal spring constant}$$

$$M_k^R = -c_\psi \psi_k \quad c_\psi \dots \text{rotational spring constant}$$

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\mathbf{B}_k^L ...unknown resulting force

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linked to the blocked degrees of freedom due to bearings!

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Bearings

We define the subsets $\beta_k^r \subset \beta_k$ and $\beta_k^h \subset \beta_k$

- Pinned

$$\mathbf{U}_{m|k} = 0 \quad m \in \beta_k$$

- Roller

$$\mathbf{U}_{m|k} \cdot \mathbf{e}_\eta = 0 \quad m \in \beta_k$$

- Clamped bearing

$$\mathbf{U}_{m|k} = 0 \quad m \in \beta_k$$

$$\psi_{m|k} = 0 \quad m \in \beta_k^r$$

- Parallel guide

$$\mathbf{U}_{m|k} \cdot \mathbf{e}_\eta = 0 \quad m \in \beta_k$$

$$\psi_{m|k} = 0 \quad m \in \beta_k^r$$

Coupling conditions

At node k we define a fixed beam

- $m_f \in \beta_k$

$$\mathbf{U}_{m_f|k} = \mathbf{U}_{m|k} \quad \forall m \in \beta_k \setminus m_f$$

- $m_r \in \beta_k^r$

$$\psi_{m_r|k} = \psi_{m|k} \quad \forall m \in \beta_k^r \setminus m_r$$

- $m_h \in \beta_k^h$

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- The governing differential equation $W^{IV}(x) + dW''(x) + eW(x) = 0$
 - Timoshenko beam theory
 - undamped, free vibration
 - constant cross section
 - frequency domain
- with the abbreviations

$$d = a + b + c \quad e = ab$$

$$a = \frac{\omega^2 \rho}{kG} \quad b = \frac{\omega^2 \rho}{E} - c \quad c = \frac{GkA}{EI}$$

- The characteristic polynomial

$$r^4 + dr^2 + e = 0$$

with the condition $\tilde{z} = r^2$ leads to

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■ Roots

$$\tilde{z}_1 = \frac{1}{2}(-d + \sqrt{\Delta}) \quad \tilde{z}_1 \in \mathbb{R} \quad \forall \omega$$

$$\tilde{z}_2 = \frac{1}{2}(-d - \sqrt{\Delta}) \quad \tilde{z}_2 \in \mathbb{R}^{<0} \quad \forall \omega$$

where

$$\Delta = d^2 - 4e \quad \Delta \wedge d \in \mathbb{R}^{>0} \quad \forall \omega$$

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- First spectrum of natural frequencies

$$\tilde{z}_1 \in \mathbb{R}^{>0} \ni \omega < \omega_{co}$$

- General solution

$$W_m(\xi) = C_{m,1} e^{-i\sqrt{\tilde{z}_2}\xi} + C_{m,2} e^{i\sqrt{\tilde{z}_2}\xi} + C_{m,3} e^{-\sqrt{\tilde{z}_1}\xi} + C_{m,4} e^{\sqrt{\tilde{z}_1}\xi}$$

- Common homogeneous solution

$$W_m(\xi) = P_{m,1} \cos(\lambda_2 \xi) + P_{m,2} \sin(\lambda_2 \xi) + P_{m,3} \cosh(\lambda_1 \xi) + P_{m,4} \sinh(\lambda_1 \xi) \quad (1)$$

where

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$$\cos(\xi) + \sin(\xi) = \left(\frac{1}{2} + \frac{i}{2} \right) e^{-i\xi} + \left(\frac{1}{2} - \frac{i}{2} \right) e^{i\xi}$$

- Implemented in the general solution

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- Dividing last term by $e^{\lambda_1 L}$ so that

$$e^{\lambda_1 \xi} e^{-\lambda_1 L} = e^{\lambda_1 (\xi - L)}$$

- Improved homogeneous solution

$$W_m(\xi) = P_{m,1} \cos(\lambda_2 \xi) + P_{m,2} \sin(\lambda_2 \xi) + C_{m,3} e^{-\lambda_1 \xi} + C_{m,4}^* e^{\lambda_1 (\xi - L)} \quad (3)$$

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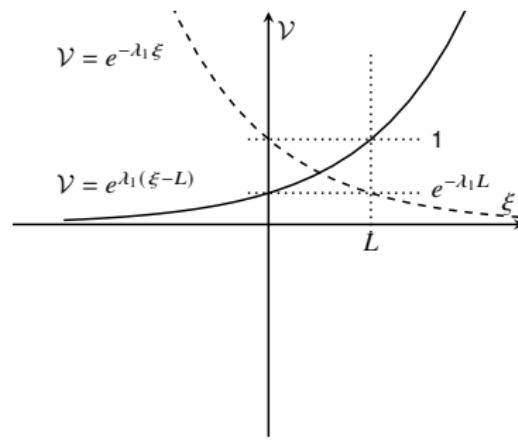
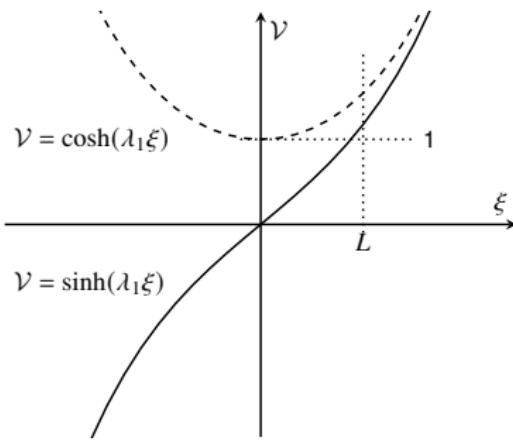
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Study of function terms

- The value range of each term in eq. (3) is $\mathcal{V} \in [e^{-\lambda_1 L}, 1]$



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1 Homogeneous solutions

$$U_m(\xi) = Q_{m,1} \cos(c_m \xi) + Q_{m,2} \sin(c_m \xi)$$

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2 Enforcement of boundary and coupling conditions

3 System of equations

$$\mathbf{A}(\omega) \mathbf{c} = \mathbf{0} \quad \mathbf{A} \in \mathbb{R}^{n \times n} \quad n = |\mathcal{T}| \cdot 6 + |\mathcal{B}|$$

4 Non-trivial solutions for ω (Regula-Falsi Method)

$$f(\omega) = \det(\mathbf{A}) = 0$$

5 Solving for the integration constants, bearing forces and moments

6 Computing eigenvectors, mode shape

1 Homogeneous solutions

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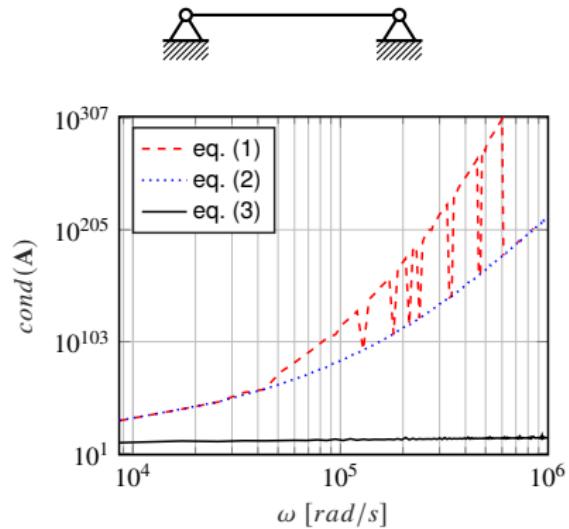
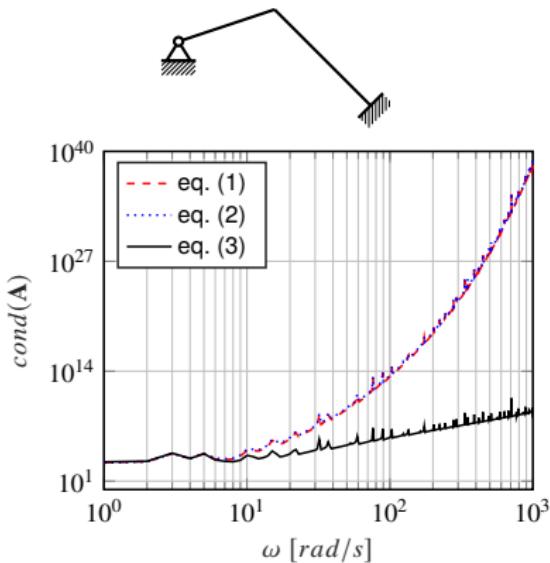
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Conditioning number



$$\text{eq. (1)} \quad W_m(\xi) = P_{m,1} \cos(\lambda_2 \xi) + P_{m,2} \sin(\lambda_2 \xi) + P_{m,3} \cosh(\lambda_1 \xi) + P_{m,4} \sinh(\lambda_1 \xi)$$

$$\text{eq. (2)} \quad W_m(\xi) = P_{m,1} \cos(\lambda_2 \xi) + P_{m,2} \sin(\lambda_2 \xi) + C_{m,3} e^{-\lambda_1 \xi} + C_{m,4} e^{\lambda_1 \xi}$$

$$\text{eq. (3)} \quad W_m(\xi) = P_{m,1} \cos(\lambda_2 \xi) + P_{m,2} \sin(\lambda_2 \xi) + C_{m,3} e^{-\lambda_1 \xi} + C_{m,4}^* e^{\lambda_1 (\xi - L)}$$

Natural frequencies

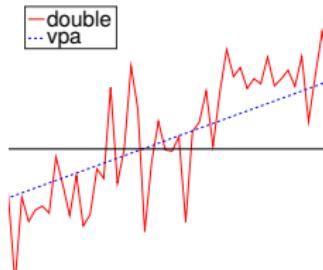
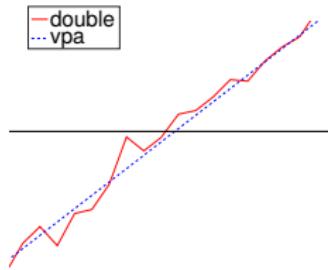
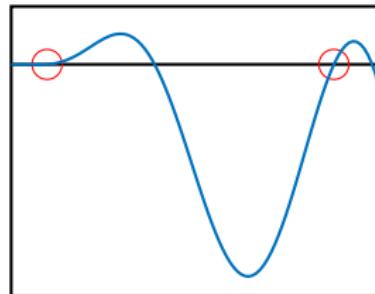
- For the arbitrary frame structure
- # $\omega_{co} = 5988$

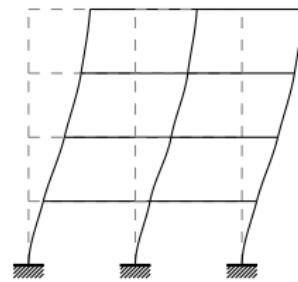
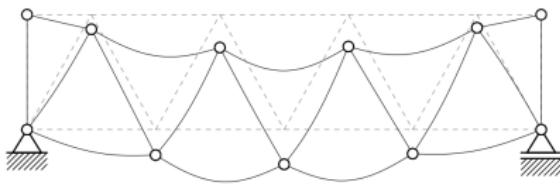
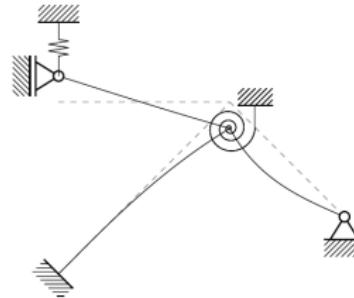
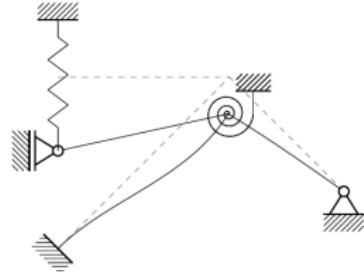
	Numerical Assembly Technique						Finite Element Method		
	double			vpa			Elements per Beam		
	(1)	(2)	(3)	(1)	(2)	(3)	16	32	64
max # ω	16	203	ω_{co}	52	273	ω_{co}	350 (25 ¹)	702 (67 ¹)	1406 (131 ¹)

$$^1 \# \omega \left| \frac{\omega - \omega_A}{\omega_A} \right| < \frac{1}{10}$$

- Zoomed-in pictures of the frequency function $f(\omega)$ at the natural frequencies ω_i
- $f(x_i)$ is plotted with $x_i \in [\omega_i - \epsilon, \omega_i + \epsilon]$ | $\epsilon = 10^{-12}$

Frequency function





Conclusion

- + "Exact" values of the natural frequencies
- + Stable computations for high frequency ranges
- + Discretization at nodes

- $f(\omega) = \det(\mathbf{A}) \rightarrow$ expensive
- Finding zeros of the frequency function in numerical matter

- Expanding to 3-D including torsional vibrations
- Implementing arches
- Considering structural damping

Further application

Automatic generation of a large number of homework assignments for basic courses in Statics and Strength of Materials

The Numerical Assembly Technique for arbitrary planar beam systems based on an improved homogeneous solution

Thomas Kramer, Michael Helmut Gfrerer

GAMM Annual Meeting
Aachen, 15. August 2022